### Almost Quaternionic Structures on Quaternionic Kaehler Manifolds

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#### Abstract

In this work, we consider a Riemannian manifold M with an almost quaternionic structure V defined by a three-dimensional subbundle of (1,1) tensors F, G and H such that  $\{F, G, H\}$  is chosen to be a local basis for V. For such a manifold there exits a subbundle  $\mathcal{H}(M)$  of the bundle of orthonormal frames  $\mathcal{O}(M)$ . If M admits a torsion-free connection reducible to a connection in  $\mathcal{H}(M)$ , then we give a condition such that the torsion tensor of the bundle vanishes. We also prove that if M admits a torsion-free connection reducible to a connection in  $\mathcal{H}(M)$ , then the tensors  $\tilde{F}^2$ ,  $\tilde{G}^2$ ,  $\tilde{H}^2$  are torsion-free, that is, they are integrable. Here  $\tilde{F}$ ,  $\tilde{G}$ ,  $\tilde{H}$  are the extended tensors of F, G and H defined on M. Finally, we show that if the torsion sof  $\tilde{F}^2$ ,  $\tilde{G}^2$  and  $\tilde{H}^2$  vanish, then Madmits a connection with torsion which is reducible to  $\mathcal{H}(M)$ , and this means that  $\tilde{F}^2$ ,  $\tilde{G}^2$  and  $\tilde{H}^2$  are integrable.

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# 1 Introduction

By imposing special structures on the tangent bundle of a manifold, it is possible to have different type of geometries. (Almost) complex, (almost) quaternionic and polynomial structures are examples of such type of structures [1-9].

In literature almost complex and almost quaternionic structures have been investigated widely, and a detailed review can be found in Kirichenko and Arseneva [5]. We begin by recalling basic results and definitions from literature.

An almost complex structure on a manifold M is a tensor field  $J: TM \to TM$ satisfying the identity  $J^2 = -id$ . An almost hypercomplex structure on a 4mdimensional manifold M is a triple S = (F, G, H) of almost complex structures F, G and H satisfying the conditions

$$F^{2} = G^{2} = H^{2} = -I, \quad H = FG, \quad FG + GF = FH + HF = GH + HG = 0,$$
(1.1)

where I denotes the identity transformation of  $T_x(M)$ . If each of the tensor fields F, G and H is a complex structure, then S is called hypercomplex structure on M.

Let M be a 4m-dimensional Riemannian manifold admitting a three dimensional subbundle V of (1,1) tensors such that on a neighborhood U of each  $x \in M$ , V has a local base  $\{F, G, H\}$ . If on each such neigborhood, the tensors F, G and H satisfy the conditions (1.1), then the bundle V is called an almost quaternionic structure on M [2].

In [10], the Nijenhuis bracket of two tensor fields A and B of type (1,1) is defined as the following tensor field of type (1,2)

$$[A, B](X, Y) = [AX, BY] - A[BX, Y] - B[X, AY] + [BX, AY] - B[AX, Y] - A[X, BY] + (AB + BA)[X, Y].$$
(1.2)

In particular, if A = B we have

$$[A, A](X, Y) = 2([AX, AY] + A^{2}[X, Y] - A[AX, Y] - A[X, AY]).$$
(1.3)

In [6], using (1.2), the torsion tensor of V is defined by [V, V] = [F, F] + [G, G] + [H, H], where [, ] denotes the Nijenhuis bracket.

The Newlander-Nirenberg Theorem states that an almost complex structure is a complex structure if and only if it is integrable i.e., it has no torsion. Thus, if the tensor fields F, G and H are integrable, then the bracket of any two of them vanishes, that is, [1]

$$[F, F] = [G, G] = [H, H] = 0$$
 and  $[F, G] = [F, H] = [G, H] = 0.$  (1.4)

Also, it is shown that there exits a torsion-free connection that F, G and H are covariantly constant which means that V is a trivial bundle [1].

In literature, moreover, it is proved that if either M is a quaternionic Kaehler manifold, or if M is a complex manifold with almost complex structures, then the vanishing of torsion tensor of V is equivalent to the vanishing of all the Nijenhuis brackets of  $\{F, G, H\}$ . For a quaternionic Kaehler manifold, with a torsion-free connection  $\nabla$ , it is given a condition implying that if [V, V] = 0 then  $\nabla F = \nabla G = \nabla H = 0$ . It follows that the bundle V admits a flat connection, hence it is trivial see [6].

In [1], a reducibility condition is given for a torsion-free connection on a quaternionic manifold M to be reducible to a connection in  $\mathcal{H}(M)$ , where  $\mathcal{H}(M)$  is subbundle of the bundle of orthonormal frame  $\mathcal{O}(M)$ .

For a Riemannian manifold M of dimension 4m + n with an almost quaternionic structure V of rank 4m, it is shown that the metric g satisfies  $g(X, \phi Y) + g(\phi X, Y) = 0$  for any vector fields  $X, Y \in \mathcal{X}(M)$ , where  $\mathcal{X}(M)$  is the algebra of vector fields on M, and any section  $\phi$  of V. In [3] Doğanaksoy studied orthogonal plane fields P and Q defined by V admitting local basis  $\{F, G, H\}$  and defined projection tensors  $p: P \to T_x M$  and  $q: Q \to T_x M$ .

In this paper, by using the projection tensors p and q above, we define extended tensors  $\widetilde{F}$ ,  $\widetilde{G}$  and  $\widetilde{H}$  on M, and we obtain reducibility conditions for a torsionfree connection on M to be a connection (either torsion-free or with torsion) in  $\mathcal{H}(M)$ . The reducibility condition corresponds to the integrability of extended tensors.

# 2 Almost Quaternionic Structures

We consider a Riemannian manifold M of dimension 4m + n admitting an almost quaternionic structure V. Let  $\{F, G, H\}$  be a local basis for V on a neighborhood U of M. Since torsion tensors [F, F], [G, G], and [H, H] are locally defined objects, to obtain a global condition for the triviality of V, a tensor [V, V] of type (1, 2) is defined globally on U by [6]

$$[V,V] = [F,F] + [G,G] + [H,H].$$
(2.1)

Let  $\{F, G, H\}$  and  $\{F', G', H'\}$  be local bases for V defined on neighborhoods U and U', respectively, and assume that  $U \cap U' \neq 0$ . Since U and U' are not disjoint, then in  $U \cap U'$  we have

$$F' = a_{11}F + a_{12}G + a_{13}H, \qquad (2.2)$$

$$G' = a_{21}F + a_{22}G + a_{23}H, \qquad (2.3)$$

$$H' = a_{31}F + a_{32}G + a_{33}H \tag{2.4}$$

where  $a_{ij}$  are functions defined on  $U \cap U'$ . With the notation above, at any point  $x \in U \cap U'$ ,  $a_{ij} \in SO(3)$  [1,2].

We calculate the torsion tensor of [V, V] defined in (2.1) by taking into account local bases  $\{F, G, H\}$  and  $\{F', G', H'\}$ . Using equations (2.2)-(2.4), we also, compute torsions [G, F], [F, H], [H, F], [G, H] and [H, G] and show that the torsion tensor [V, V] of the bundle V is independent of the choice of bases, that is, [V, V] is well defined.

If a tensor f of type (1,1) on M satisfies the structure equation

$$f^3 + f = 0, (2.5)$$

then f is called an f-structure. Since the tensors F, G and H satisfy the conditions (1.1), they are f-structures on M. In [3] it is shown that any cross-section  $\phi$  of V of length 1 is f-structure on M, and F, G and H are of length 1.

Let  $\{F, G, H\}$  be a basis for V in some neighborhood U of M. In [11], it is proved that each of F, G, and H has a constant rank on U and from the conditions F = GH, G = HF, H = FG, it is seen that their ranks are all equal. Also, in [11], the rank of V is defined to be the rank of a basis element on some neighborhood U. By choosing  $q = 1 + F^2 = 1 - p$ , where 1 denotes the identity operator, it is obtained in [3] that

$$p + q = 1, \quad p^2 = p, \quad q^2 = q$$
 (2.6)

and that

$$\phi p = p \phi = \phi, \quad \phi q = q \phi = 0, \qquad (2.7)$$

for any cross-section  $\phi$  of V. This shows that p and q are complementary projection operators. Then, there exist two distributions P and Q corresponding to p and q, respectively. If the rank of V is 4m, then P is 4m-dimensional and Q is n-dimensional.

Let g' be a Riemannian metric of M. Define g to be the tensor field of degree 2 on M as

$$g(X,Y) = \begin{cases} 0, & X \in P(x), Y \in Q(x) \\ g'(X,Y), & X,Y \in Q(x) \\ g'(X,Y) + g'(FX,FY) \\ +g'(GX,GY) + g'(HX,HY), & X,Y \in P(x) \end{cases}$$
(2.8)

where  $\{F, G, H\}$  is a canonical local base for V, and  $X, Y \in T_x(M)$ . Since g' is a Riemannian, g satisfies all the conditions for a Riemannian metric [3]. For any cross-section  $\phi$  of V, from (2.8) we have

$$g(X,\phi Y) + g(\phi X, Y) = 0, \quad X, Y \in \mathcal{X}(M).$$
(2.9)

If M admits an almost quaternionic structure, then at each point x in M there is an orthonormal basis of TM, of the form

$$\{X_1,\ldots,X_m,FX_1,\ldots,FX_m,GX_1,\ldots,GX_m,HX_1,\ldots,HX_m\},\qquad(2.10)$$

and the set of all such frames at all points  $x \in M$  constitutes a subbundle, denoted by  $\mathcal{H}(M)$ , of the bundle of orthonormal frames  $\mathcal{O}(M)$ .

We consider P and Q to be the orthogonal plane fields defined by almost substructure V, and let p and q be projections  $P \to T_x M$  and  $Q \to T_x M$ , respectively [3].

Let us denote by f the tensor Jp defined on P. Since pJ = J, f satisfies the structure equation  $f^3 + f = 0$ . For any tensor J of type (1, 1) on M we define its extension as the tensor  $\widetilde{J}$  of type (1, 1) on M by setting

$$\widetilde{J} = q + Jp = q + f.$$
(2.11)

It can be seen that the extended tensor  $\widetilde{J}$  has the following properties

$$\widetilde{J}^2 = q - p \,, \tag{2.12}$$

$$\widetilde{J}^3 = q - Jp = 1 - \widetilde{J} + \widetilde{J}^2, \qquad (2.13)$$

$$\widetilde{J}^4 = 1. \tag{2.14}$$

We recall that  $\Gamma$  is an affine (linear) connection on M. A tensor field, say J, is parallel with respect to  $\Gamma$  if and only if  $\nabla J = 0$ , where  $\nabla$  is the covariant differentiation with respect to  $\Gamma$  [10].

We give the following proposition from [4].

**Proposition.** The followings are equivalent

(a)  $[\widetilde{J}^2, \widetilde{J}^2]$  vanishes,

(b) M admits a torsion-free affine connection according to which  $\tilde{J}^2$  is parallel,

(c) the plane field P and its orthogonal complement Q are both integrable.

# 3 Torsion-free connections reducible to a connection in $\mathcal{H}(M)$

Compatibility of a torsion-free connection  $\Gamma$  with the metric g' is equivalent to the reducibility of  $\Gamma$  to  $\mathcal{O}(M)$ . Furthermore, if  $\Gamma$  is reducible to  $\mathcal{H}(M)$ , then the manifold is called *quaternionic Kaehler manifold* [2].

We also quote the following theorem from [2].

**Theorem.** Let M be an almost quaternionic manifold. A connection  $\Gamma$  in orthonormal frame  $\mathcal{O}(M)$  is reducible to a connection in  $\mathcal{H}(M)$  if and only if the covariant derivatives of the tensor fields of F, G and H satisfy the following conditions:

$$\nabla F = aG - bH, \quad \nabla G = -aF + cH, \quad \nabla H = bF - cG$$
(3.1)

where a, b and c are 1-forms on M.

We now begin to state our main theorems related to the reducibility of a torsion free-connection.

**Theorem 3.1.** If M admits a torsion-free connection  $\Gamma$  reducible to  $\mathcal{H}(M)$ such that

$$(\nabla\phi)(\phi X, Y) + (\nabla\phi)(X, \phi Y) = 0 \tag{3.2}$$

for any section  $\phi \in \{F, G, H\}$ , then the torsion tensor [V, V] of the bundle V vanishes.

**Proof.** For any section  $\phi \in \{F, G, H\}$ , from (3.2) we have

$$(\nabla F)(FX,Y) + (\nabla F)(X,FY) = 0, \qquad (3.3)$$

$$(\nabla G)(GX,Y) + (\nabla G)(X,GY) = 0, \qquad (3.4)$$

$$(\nabla H)(HX,Y) + (\nabla H)(X,HY) = 0.$$
 (3.5)

By using (3.1) in (3.3)-(3.5), we obtain following relations

$$(\nabla F)(FX,Y) = (\nabla_Y F)(FX) = a(Y)G(FX) - b(Y)H(FX)$$
  
=  $-a(Y)H(X) - b(Y)G(X)$ , (3.6)

$$(\nabla F)(X, FY) = (\nabla_{FY}F)(X) = a(FY)G(X) - b(FY)H(X). \quad (3.7)$$

Similarly, we find

$$(\nabla G)(GX,Y) = (\nabla_Y G)(GX) = -a(Y)F(GX) + c(Y)H(GX)$$
  
=  $-a(Y)H(X) - c(Y)F(X)$ , (3.8)

$$(\nabla G)(X, GY) = (\nabla_{GY}G)(X) = -a(GY)F(X) + c(GY)H(X) \quad (3.9)$$

and

$$(\nabla H)(HX,Y) = (\nabla_Y H)(HX) = b(Y)F(HX) - c(Y)G(HX) = -b(Y)G(X) - c(Y)F(X), \qquad (3.10) (\nabla H)(X,HY) = (\nabla_{HY}H)(X) = b(HY)F(X) + b(Y)G(X). \qquad (3.11)$$

$$\nabla H)(X, HY) = (\nabla_{HY}H)(X) = b(HY)F(X) + b(Y)G(X)$$
. (3.11)

For a torsion-free connection  $\Gamma$  reducible to  $\mathcal{H}(M)$ , if the 1-forms a, b and c defined in (3.1) are used in the equations (3.6)-(3.11), then we observe that they satisfy the following relations

$$b(X) = a(FX),$$
  

$$c(X) = -a(GX),$$
  

$$c(X) = b(HX).$$
(3.12)

In order to find the torsion tensor of V, we compute the torsions of tensors F, G and H:

$$\frac{1}{2}[F,F](X,Y) = ([FX,FY] + F^2[X,Y] - F[FX,Y] - F[X,FY]$$

$$= (\nabla_{FX}FY - \nabla_{FY}FX) + F^2(\nabla_XY - \nabla_{FY}X)$$

$$-F(\nabla_{FX}Y - \nabla_YFX) - F(\nabla_XFY - \nabla_{FY}X)$$

$$\frac{1}{2}[F,F](X,Y) = (\nabla_{FX}F)Y - (\nabla_{FY}F)X + F((\nabla_YF)X) - F((\nabla_XF)Y).$$
(3.13)

Similarly, the tensors [G, G] and [H, H] are obtained as follows:

$$\frac{1}{2}[G,G](X,Y) = (\nabla_{GX}G)Y - (\nabla_{GY}G)X + G((\nabla_YG)X) - G((\nabla_XG)Y)$$
(3.14)

and  

$$\frac{1}{2}[H,H](X,Y) = (\nabla_{HX}H)Y - (\nabla_{HY}H)X + H((\nabla_{Y}H)X) - H((\nabla_{X}H)Y).$$
(3.15)

Using the relations of (3.12), we get the following equalities

$$\frac{1}{2}[F,F](X,Y) = [a(FX)G(Y) - b(FX)H(Y)] - [a(FY)G(X) - b(FY)H(X)] 
+ F[a(Y)G(X) - b(Y)H(X)] - F[a(X)G(Y) - b(X)H(Y)] 
\frac{1}{2}[F,F](X,Y) = (b(Y) - a(FY))G(X) + (a(FX) - b(X))G(Y) 
+ (b(FY) + a(Y))H(X) + (-a(X) - b(FX))H(Y).$$
(3.16)

By using (3.14) and (3.15) for G and H, we obtain

$$\begin{aligned} \frac{1}{2}[G,G](X,Y) &= (-a(GX) - c(X))F(Y) + (a(GY) + c(Y))F(X) \\ &+ (a(Y) - c(GY))H(X) + (c(GX) - a(X))H(Y) \end{aligned} \tag{3.17}$$

and

$$\frac{1}{2}[H,H](X,Y) = (b(HX) - c(X))F(Y) + (-b(HY) + c(Y))F(X) + (b(Y) + c(HY))G(X) + (-b(X) - c(HX))G(Y).$$
(3.18)

Substituting the conditions (3.12) in (3.16)-(3.18), we find [F, F] = 0, [G, G] = 0, [H, H] = 0, and therefore [V, V] = 0. This completes the proof.

Given tensor fields F, G and H, let us define their extended tensors  $\widetilde{F}$ ,  $\widetilde{G}$  and  $\widetilde{H}$  as in (2.11). It is clear that the relations (2.12)-(2.14) hold for the tensors  $\widetilde{F}$ ,  $\widetilde{G}$  and  $\widetilde{H}$ .

To simplify our calculations we use the following notations

$$\widetilde{F} = q + Fp = q + f, \qquad (3.19)$$

$$\widetilde{G} = q + Gp = q + g, \qquad (3.20)$$

$$\widetilde{H} = q + Hp = q + h, \qquad (3.21)$$

where f = Fp, g = Gp, and h = Hp. Since p + q = 1, where 1 denotes the identity operator, and tensors F, G, H are defined on distribution P of TM, using the projection tensor p of type (1,1) defined in Section 2, we write  $q = 1 + f^2 = 1 - p$ . Then, we have

$$p^2 = p, \quad q^2 = q.$$
 (3.22)

Also using (2.12), (2.13), and (2.14), we show that the following relations hold:

$$\widetilde{F}^2 = q - p \,, \tag{3.23}$$

$$\widetilde{F}^{3} = \widetilde{F}\widetilde{F}^{2} = (q+f)(q-p) = q-f$$
, (3.24)

$$\widetilde{F}^3 + \widetilde{F} = 2q. aga{3.25}$$

Similar calculations can be done for  $\widetilde{G}$  and  $\widetilde{H}$  by letting  $q = 1 + g^2 = 1 - p$  and  $q = 1 + h^2 = 1 - p$ , respectively.

**Theorem 3.2.** If M admits a torsion-free connection  $\Gamma$  reducible to  $\mathcal{H}(M)$ , then the tensors  $\widetilde{F}^2 = \widetilde{G}^2 = \widetilde{H}^2 = q - p$  are torsion-free. That is,  $[\widetilde{F}^2, \widetilde{F}^2] = 0, \ [\widetilde{G}^2, \widetilde{G}^2] = 0$  and  $[\widetilde{H}^2, \widetilde{H}^2] = 0$ , i.e., P and Q are integrable. **Proof.** Since  $\Gamma$  is reducible to  $\mathcal{H}(M)$ , equations in (3.1) hold. We get  $\nabla p = -\nabla q$  from p + q = 1. Since  $q = 1 + f^2 = 1 - p$ , using (3.22) we obtain

$$\nabla p = -\nabla(f^2) = -\nabla(Fp \circ Fp) = -Fp\nabla(Fp) - \nabla(Fp)Fp.$$
(3.26)

By using (3.1) we get

$$\nabla p = -[(aG - bH)Fp + F(\nabla p)Fp + Fp(aG - bH)p - (\nabla p)], \qquad (3.27)$$

which implies

$$\nabla p = 0. \tag{3.28}$$

So,  $\nabla p$  and  $\nabla q$  are both 0. It means that p and q are parallel. Furthermore,  $\widetilde{F}^2 = q - p$  implies that  $\nabla \widetilde{F}^2 = 0$ . Then,  $[\widetilde{F}^2, \widetilde{F}^2] = 0$ . Similarly, we see that  $[\widetilde{G}^2, \widetilde{G}^2] = 0$  and  $[\widetilde{H}^2, \widetilde{H}^2] = 0$ . Hence, P and Q are integrable by Proposition.

**Theorem 3.3.** If the torsions of  $\widetilde{F}^2$ ,  $\widetilde{G}^2$ , and  $\widetilde{H}^2$  vanish with respect to a torsion-free connection  $\Gamma$ , then M admits a connection  $\overset{\circ}{\Gamma}$  reducible to  $\mathcal{H}(M)$  whose torsion tensor is given by

$$24 T(X,Y) = 2 \left( [f,f] + [g,g] + [h,h] \right) (X,Y) - \left( [f,f] + [g,g] + [h,h] \right) (pX,pY)$$
(3.29)

for any vector fields X and Y on M.

**Proof.** Since  $\Gamma$  is an arbitrary torsion-free affine connection on M with covariant derivative  $\nabla$ , by using Proposition, we have  $\nabla \tilde{F}^2 = 0$ ,  $\nabla \tilde{G}^2 = 0$  and  $\nabla \tilde{H}^2 = 0$ .

Below we give calculations only for the tensor  $\widetilde{F}$ . The calculations for the other tensors  $\widetilde{G}$  and  $\widetilde{H}$  can be obtained similarly.

To simplify the calculations, we introduce the tensors  $W_I$  and  $W_{II}$  of type (1, 2) as follows:

$$W_{I}(X,Y) = (\nabla_{\widetilde{F}Y}\widetilde{F})X + \widetilde{F}((\nabla_{Y}\widetilde{F})X) + \widetilde{F}^{2}((\nabla_{\widetilde{F}^{3}Y}\widetilde{F})X) + \widetilde{F}^{3}((\nabla_{\widetilde{F}^{2}Y}\widetilde{F})X),$$
(3.30)
$$W_{II}(X,Y) = \widetilde{F}((\nabla_{X}\widetilde{F})Y)$$
(3.31)

for any vector fields X and Y on M. Using (3.30) we obtain

$$W_{I}(X, \widetilde{F}Y) = (\nabla_{\widetilde{F}^{2}Y}\widetilde{F})X + \widetilde{F}((\nabla_{\widetilde{F}Y}\widetilde{F})X) + \widetilde{F}^{2}((\nabla_{\widetilde{Y}}\widetilde{F})X) + \widetilde{F}^{3}((\nabla_{\widetilde{F}^{3}Y}\widetilde{F})X)$$

$$(3.32)$$

and

$$\widetilde{F}W_{I}(X,Y) = \widetilde{F}((\nabla_{\widetilde{F}Y}\widetilde{F})X) + \widetilde{F}^{2}((\nabla_{Y}\widetilde{F})X) + \widetilde{F}^{3}((\nabla_{\widetilde{F}^{3}Y}\widetilde{F})X) + ((\nabla_{\widetilde{F}^{2}Y}\widetilde{F})X),$$
(3.33)

which gives

$$W_I(X, \widetilde{F}Y) - \widetilde{F}(W_I(X, Y)) = 0.$$
(3.34)

From (3.31) we also obtain

$$W_{II}(X, \widetilde{F}Y) - \widetilde{F}(W_{II}(X, Y)) = 2(\nabla_X \widetilde{F})Y.$$
(3.35)

As the tensor  $\widetilde{F}^2$  is parallel, i.e,  $\nabla \widetilde{F}^2 = 0$ , we have  $(\nabla \widetilde{F})\widetilde{F} = -\widetilde{F}(\nabla \widetilde{F})$ . Using (2.6), (2.7), and (2.8), we see that the following relations hold:

$$\nabla \widetilde{F} = \nabla q + \nabla F p = \nabla F p = \nabla f \,. \tag{3.36}$$

Since  $\nabla \widetilde{F}^2 = 0$ ,  $\nabla p = \nabla q = 0$ . Thus,  $\nabla \widetilde{F}^3 = -\nabla \widetilde{F}$ .

For any tensor field  $\alpha$ , if  $\alpha$  is parallel, then  $\nabla_X \alpha Y = \alpha(\nabla_X Y)$ . Therefore,

$$\nabla \widetilde{F}^3 = \nabla (\widetilde{F}^2 \circ \widetilde{F}) = \nabla \widetilde{F}^2 (\nabla \widetilde{F}) \,. \tag{3.37}$$

In addition to the equations (3.34) and (3.35), the following relations hold:

$$W_{II}(X, \widetilde{F}Y) = \widetilde{F}((\nabla_X \widetilde{F})\widetilde{F}Y) = -\widetilde{F}^2((\nabla_X \widetilde{F})Y), \qquad (3.38)$$

$$\widetilde{F}W_{II}(X,Y) = \widetilde{F}^2((\nabla_X \widetilde{F})Y) = (\nabla_X \widetilde{F}^3)Y = (\nabla_X \widetilde{F})Y.$$
(3.39)

From (3.38) and (3.39) we get

$$W_{II}(X,\widetilde{F}Y) - \widetilde{F}(W_{II}(X,Y)) = -2(\nabla_X \widetilde{F}^3)Y = 2(\nabla_X \widetilde{F})Y. \quad (3.40)$$

Rearranging (3.30) we have

$$W_{I}(X,Y) = (\nabla_{qY+fY}\widetilde{F})X + \widetilde{F}((\nabla_{pY+qY}\widetilde{F})X) + \widetilde{F}^{2}((\nabla_{qY-FY}\widetilde{F})X) + \widetilde{F}^{3}((\nabla_{qY-pY}\widetilde{F})X) = ((\nabla_{qY}\widetilde{F})X)(1 + \widetilde{F} + \widetilde{F}^{2} + \widetilde{F}^{3}) + ((\nabla_{FY}\widetilde{F})X)(1 - \widetilde{F}^{2}) + ((\nabla_{pY}\widetilde{F})X)(\widetilde{F} - \widetilde{F}^{3}) = 4q((\nabla_{qY}\widetilde{F})X) + 2p((\nabla_{fY})\widetilde{F})X) + 2f((\nabla_{pY}\widetilde{F})X).$$
(3.41)

Because  $\nabla p = \nabla q = 0$ , we obtain  $\nabla \tilde{F} = \nabla F p$ . So, by the definition of the projection tensor q, we obtain  $q(\nabla \tilde{F}) = 0$ . Moreover,

$$q(\nabla_X F)Y = 0$$
 for any vector fields X, and  $Y \in T_x(M)$ , (3.42)

and

$$q(\nabla_X F)Y = q((\nabla_X (FY) - F(\nabla_X Y)))$$
  
=  $q(\nabla_X (FY)) - qF(\nabla_X Y) = q(\nabla_X (FY)) = 0.$  (3.43)

Since q is parallel, it follows that  $q(\nabla_X(FY)) = \nabla_X(qFY) = 0$ , which reduces (3.41) to

$$W_{I}(X,Y) = 4q[(\nabla_{qY}q)(X) + (\nabla_{qY}f)X] + 2p[(\nabla_{fY}q)X + (\nabla_{fY}f)X] + 2f[(\nabla_{pY}q)X + (\nabla_{pY}f)X].$$
(3.44)

Furthermore, we have  $q((\nabla_{qY}(FX)) = p((\nabla_{FY}q)X) = F((\nabla_{pY}q)X) = 0$ , which together with (3.43), implies

$$W_I(X,Y) = 2p[(\nabla_{fY}f)X] + 2f[(\nabla_{pY}f)X].$$
(3.45)

(3.31) and (3.19) give

$$W_{II}(X,Y) = \widetilde{F}((\nabla_X \widetilde{F})Y)$$
  
=  $(q+f)(\nabla_X (q+f)Y)$   
=  $q((\nabla_X q)Y) + q((\nabla_X f)Y) + f((\nabla_X q)Y) + f((\nabla_X f)Y),$   
 $W_{II}(X,Y) = f((\nabla_X f)Y).$  (3.46)

Let us now consider an affine connection  $\overset{\circ}{\Gamma}$  whose covariant derivative  $\overset{\circ}{\nabla}$  is given by

$$\overset{\circ}{\nabla}_{X}Y = \nabla_{X}Y - \frac{1}{8}W_{I}(X,Y) - \frac{1}{2}W_{II}(X,Y), \qquad (3.47)$$

where X and Y are vector fields on M. Here  $\nabla$  is covariant differentiation of an affine connection, and  $W_I$  and  $W_{II}$  are bilinear mappings  $T_x(M) \times T_x(M) \rightarrow$  $T_x(M)$  at each  $x \in M$ , and  $\stackrel{\circ}{\nabla}$  defines an affine connection on M.

We can see that by using (3.47),  $\widetilde{F}$ ,  $\widetilde{G}$ , and  $\widetilde{H}$  are parallel with respect to  $\overset{\circ}{\Gamma}$  on M.

Let X and Y be arbitrary vector fields on M. Then,

$$\begin{aligned} (\mathring{\nabla}_{X}\widetilde{F})Y &= \mathring{\nabla}_{X}(\widetilde{F}Y) - \widetilde{F}(\mathring{\nabla}_{X}Y) \\ &= \nabla_{X}(\widetilde{F}Y) - \frac{1}{8}W_{I}(X,\widetilde{F}Y) - \frac{1}{2}W_{II}(X,\widetilde{F}Y) \\ &-\widetilde{F}[\nabla_{X}Y - \frac{1}{8}W_{I}(X,Y) - \frac{1}{2}W_{II}(X,Y)] \\ &= \nabla_{X}(\widetilde{F}Y) - \widetilde{F}(\nabla_{X}Y) - \frac{1}{8}(W_{I}(X,\widetilde{F}Y) - F(W_{II}(X,\widetilde{F}Y))) \\ &-\frac{1}{2}[W_{II}(X,\widetilde{F}Y) - \widetilde{F}(W_{II}(X,Y))] \\ &= (\nabla_{X}\widetilde{F})Y - \frac{1}{2}(2(\nabla_{X}\widetilde{F})Y) = 0, \end{aligned}$$
(3.48)

hence we conclude that  $\overset{\circ}{\nabla}_{X}\widetilde{F} = 0$ . By (3.47), the torsion of  $\overset{\circ}{\Gamma}$  is obtained as

$$T(X,Y) = \overset{\circ}{\nabla}_{X}Y - \overset{\circ}{\nabla}_{Y}X - [X,Y]$$
  
=  $\nabla_{X}Y - \nabla_{Y}X - [X,Y] - \frac{1}{8}(W_{I}(X,Y) - W_{I}(Y,X))$   
 $-\frac{1}{2}(W_{II}(X,Y) - W_{II}(Y,X)).$  (3.49)

Since  $\Gamma$  is torsion-free , we get

$$T(X,Y) = -\frac{1}{8}(W_I(X,Y) - W_I(Y,X)) - \frac{1}{2}(W_{II}(X,Y) - W_{II}(Y,X)). \quad (3.50)$$

Using (3.45) and (3.46) we find

$$-8T(X,Y) = W_{I}(X,Y) - W_{I}(Y,X) + 4W_{II}(X,Y) - 4W_{II}(Y,X)$$
  
=  $2p((\nabla_{fY}f)X) + 2f((\nabla_{pY}f)X) - 2p((\nabla_{fX}f)Y)$   
 $-2f((\nabla_{pX}f)Y) + 4f((\nabla_{X}f)Y) - 4f((\nabla_{Y}f)X).$  (3.51)

Because  $\nabla \widetilde{F}^2 = 0$  and  $\nabla p = 0$ , we see that  $p(\nabla Fp) = p(\nabla f) = \nabla f$ . Then, (3.51) is reduced to

$$-8T(X,Y) = 2(\nabla_{fY}f)X + 2f((\nabla_{pY}f)X) - 2(\nabla_{fX}f)Y - 2f((\nabla_{pX}f)Y) +2f((\nabla_{X})Y) + 2f((\nabla_{pX}f)Y) + 2(\nabla_{qX}f)Y - 2f((\nabla_{Y}f)X) -2f((\nabla_{pY}f)X) - 2f((\nabla_{pY}f)X) - 2f((\nabla_{qY}f)X).$$
(3.52)

Equation (3.52) can be simplified to

$$-8T(X,Y) = -2[(\nabla_{fX}f)Y - (\nabla_{fY}f)X - f((\nabla_Xf)Y) + f((\nabla_Yf)X)] +2f((\nabla_{qX}f)Y) - 2f((\nabla_{qY}f)X).$$
(3.53)

By means of (1.3), we get

$$\frac{1}{2}[f,f](X,Y) = (\nabla_{fX}f)Y - (\nabla_{fY}f)X - f((\nabla_Xf)Y) + f((\nabla_Yf)X). \quad (3.54)$$

Substituting (3.54) in (3.53) yields

$$8T(X,Y) = [f,f](X,Y) - 2f((\nabla_{qX}f)Y) + 2F((\nabla_{qY}f)X).$$
(3.55)

If torsion tensor of [f, f] is calculated for vector fields pX and pY, by using (3.54), we get

$$[f, f](pX, pY) - [f, f](X, Y) = 2([fX, fY] - p[pX, pY] - f[pX, fY] - f[fX, pY] - ([fX, fY] + f2[X, Y] - f[fX, Y] - f[X, fY])).$$
(3.56)

Moreover, we have

$$f[fX,Y] = f[fX,pY+qY] = f[fX,pY] + f[fX,qY],$$
  

$$f[X,fY] = f[pX,fY] + f[qX,fY].$$
(3.57)

Since  $(Fp) \circ (Fp) = -p$  and  $f^2 = -p$ , we obtain

$$f^{2}[X,Y] = -p[pX + qX, pY + qY] = -p[pX, pY] - p[pX, qY] - p[qX, pY].$$
(3.58)

Substituting (3.57) and (3.58) in (3.56), we get

$$[f, f](pX, pY) - [f, f](X, Y) = 2(p[qX, pY] + p[pX, qY] + p[qX, qY] + f[fX, qY] + f[qX, fY]).$$
(3.59)

As  $\nabla \widetilde{F}^2=0$  and  $\Gamma$  is torsion-free, by Proposition we have p[qX,qY]=0. So , (3.59) becomes

$$[f, f](pX, pY) - [f, f](X, Y) = 2(p(\nabla_{pX}qY - \nabla_{qY}pX) + p(\nabla_{qX}pY - \nabla_{pY}qX) + f(\nabla_{fX}qY - \nabla_{qY}fX) + f(\nabla_{qX}fY) - \nabla_{fY}qX).$$
(3.60)

By means of parallel properties of p and q, we get

$$p(\nabla_{pY}(qX)) = p(\nabla_{pX}(qY)) = F(\nabla_{FY}(qX)) = 0$$
(3.61)

and

$$p(\nabla_{qX}(pY)) = p(\nabla_{qX}Y), \quad p(\nabla_{qY}(pX)) = p(\nabla_{qY}X).$$

Therefore, (3.60) reduces to

$$[f, f](pX, pY) - [f, f](X, Y) = 2f((\nabla_{qX} f)Y) - 2f((\nabla_{qY} f)X).$$

Finally, torsion of f becomes

$$8T(X,Y) = 2[f,f](X,Y) - [f,f](pX,pY).$$
(3.62)

Similarly, torsion tensors for g and h are obtained as

$$8T(X,Y) = 2[g,g](X,Y) - [g,g](pX,pY), \qquad (3.63)$$

$$8T(X,Y) = 2[h,h](X,Y) - [h,h](pX,pY), \qquad (3.64)$$

which completes the proof.

# References

- Yano, K., Ako, M., "Integrability conditions for almost quaternion structure". Hokaido Math J. 1, 63-86 (1972).
- [2] Ishihara, S., "Quaternion Kahlerian Manifolds". J. Differential Geometry 9, 483-500 (1974).
- [3] Doğanaksoy, A., "Almost Quaternionic Substructures". Turkish Jour. of Math. 16, 109-118 (1992).
- [4] Doğanaksoy, A., "On Plane Fields With An Almost Complex Structure". Turkish Jour. of Math. 17, 11-17.(1993)
- [5] Kirichenko, V.F., Arseneva, O.E., "Differential geometry of generalized almost quaternionic structures. I", dg-ga/9702013.
- [6] Ozdemir, F., "A Global Condition for the Triviality of an almost quaternionic structure on complex manifolds". Int. Journal of Pure and Applied Math. Sciences 3, No.1, 1-9 (2006).
- [7] Özdemir, F., Crasmareanu, M., "Geometrical objects associated to a substructure". Turkish Jour. of Math. 35, 717-728, (2011).
- [8] Ata, E. and Yaylı Y., "A global condition for the triviality of an almost split quaternionic structure on split complex manifolds". Int. J. Math. Sci. (WASET) 2, no.1, 47-51 (2008).
- [9] Alagöz, Y., Oral, K.H., Yüce, S., "Split Quaternion Matrices". Accepted for publication in Miskolc Mathematical Notes.

- [10] Kobayashi, S., Nomizu, K., "Foundations of Differential Geometry". A Willey-Interscience Publication, Vol.1,2 (1969).
- [11] Stong, R.E., "The Rank of an  $f\mathchar`$  Kodai Math. Sem. Rep. 29, 207-209, (1977).