# Qualitative properties and standard estimates of solutions for some fourth order elliptic systems

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**Abstract:** In this paper, first, we make the uniform estimates for a class of fourth order elliptic system in bounded and smooth domains. Second, we study the qualitative properties of solutions with prescribed integration in  $\mathbb{R}^4$ . Finally, we also will obtain some radially symmetric results by using moving planes methods.

**Keywords:** Fourth order elliptic systems; Asymptotic behavior; Uniform priori estimates; Q curvature

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#### 1. Introduction

In this paper, we make uniform estimates to the following fourth order elliptic system:

$$\begin{cases} \Delta^2 u = Q_1(x)e^{4v}, & \text{in } \Omega \subset R^4; \\ \Delta^2 v = Q_2(x)e^{4u}, & \text{in } \Omega \subset R^4; \\ u = \Delta u = v = \Delta v = 0 & \text{on } \partial\Omega \end{cases}$$
(\*)

and investigate properties of the solutions to the following fourth order elliptic system:

$$\begin{cases} \Delta^2 u = Q_1(x)e^{4v}, & x \in R^4; \\ \Delta^2 v = Q_2(x)e^{4u}, & x \in R^4. \end{cases}$$
(\*\*)

Where  $\Omega$  is a bounded smooth domain in  $R^4$  and  $\Delta^2$  is the biharmonic operator.  $Q_i(x)$  (i = 1, 2) are given function in  $L^p(\Omega)$  for some 1 . We assume that $<math>u, v \in L^1(\Omega), e^{4u}, e^{4v} \in L^{p'}(\Omega)$  (where p' is the conjugate exponent of p) so that (\*) has a meaning in the sense of distributions. A first question is whether one can conclude that all eventual solutions of (\*) have uniform bounds. As we will see in the Section 2 (see our main result, Theorem 2.4) the answer is positive. It is obvious that we have generalized part works in [6] when u = v in (\*).

Recently, a series of works have been done to understand the existence and the qualitative properties of the solutions of (\*\*) when  $Q_1 = Q_2$  and u = v. In the mean time, (\*\*) becomes

$$\Delta^2 u = Q(x)e^{4u}, \quad x \in \mathbb{R}^4, \tag{(***)}$$

where  $Q = Q_1 = Q_2$ .

When Q = 6, Lin [1] had given a complete classification of u in terms of its growth, or of the behavior of  $\Delta u$  at  $\infty$ . Xu [2] had done similar work by using moving spheres methods. Wei and Xu [4] and Martinazzi [7] also gave a complete classification of solutions for higher order conformally invariant equations compared to (\*\*\*). In Section 3, we consider more general functions  $Q_1(x), Q_2(x)$  for problem (\*\*). This is considered as the generalization of problem (\* \*\*) in [1]. First, we obtain the asymptotic behavior of solutions near infinity. Consequently, we prove that all solutions satisfy an identity, which is similar to the well-known Kazdan-Warner condition (see our main result, Theorem 3.2). Finally, using the harmonic asymptotic expansion at  $\infty$  in [1], we show that all the solutions for (\*\*) are radial symmetric provided  $Q_1, Q_2$  is radially symmetric and non-increasing. This part can be viewed as the completion of [1].

#### 2. Uniform estimates for problem (\*)

Assume  $\Omega \subset \mathbb{R}^4$  is a bounded domain and let h be a solution of

$$\begin{cases} \Delta^2 h(x) = f(x), & \text{in } \Omega \subset R^4; \\ h = \Delta h = 0 & \text{on } \partial \Omega. \end{cases}$$
(2.1)

Following the argument Brezis-Merle [8], Lin obtained the following Lemma:

**Lemma 2.1.** [1] Suppose  $f \in L^1(\overline{\Omega})$ . For any  $\delta \in (0, 32\pi^2)$ , there exists a constant  $C_{\delta} > 0$  such that the inequality,

$$\int_{\Omega} \exp(\frac{\delta |h|}{||f||_{L^1}}) dx \le C_{\delta} (diam\Omega)^4,$$

where  $diam\Omega$  denotes the diameter of  $\Omega$ .

By using above Lemma, we obtain following consequent results:

**Theorem 2.1.** Let u be a solution of (2.1) with  $f \in L^1(\Omega)$ . Then for every constant k > 0,

$$e^{ku} \in L^1(\Omega).$$

**Proof.** Let  $0 < \epsilon < \frac{1}{k}$ , we may split f as  $f = f_1 + f_2$  with  $||f_1||_1 < \epsilon$  and  $f_2 \in L^{\infty}(\Omega)$ . Write  $u_i$  are the solutions of

$$\begin{cases} \Delta^2 u_i = f_i, & \text{in } \Omega; \\ u_i = \Delta u_i = 0 & \text{on } \partial \Omega. \end{cases}$$

By Lemma 2.1, we find  $\int_{\Omega} \exp[\frac{|u_1(x)|}{||f_1||_1}] < \infty$  and thus  $\int_{\Omega} \exp[k|u_1|] < \infty$ . The conclusion follows since  $|u| \le |u_1| + |u_2|$  and  $u_2 \in L^{\infty}(\Omega)$ .

Before stating our main results on a priori bounds, we state a result on the regularity of the distribution solutions of (\*\*).

**Theorem 2.2.** Suppose (u, v) is a solution of equation (\*\*) with  $Q_1, Q_2 \in L^p(\Omega)$ and  $e^{4u}, e^{4v} \in L^{p'}(\Omega)$  for some  $1 . Then <math>u, v \in L^{\infty}(\Omega)$ .

**Proof.** By Theorem 2.1, we know that  $e^{ku} \in L^1(\Omega)$  for all k, i.e.,  $e^u \in L^r(\Omega) \ \forall r < \infty$ . It follows that  $Q_2 e^{4u} \in L^{p-\delta} \ \forall \delta > 0$  if  $p < \infty$ , and  $Q_2 e^{4u} \in L^r(\Omega) \ \forall r < \infty$  if  $p = \infty$ . Standard elliptic estimates imply that  $\Delta v \in L^{\infty}(\Omega)$ . Hence, combing v = 0 on  $\partial\Omega$ , we have  $v \in L^{\infty}(\Omega)$ . Similarly, we have  $u \in L^{\infty}(\Omega)$ .

**Corollary 2.1.** Suppose (u, v) is a solution of

$$\begin{cases} \Delta^2 u = Q_1 e^{4v} + f(x), & \text{ in } \Omega; \\ \Delta^2 v = Q_2 e^{4u} + g(x), & \text{ in } \Omega; \\ u = g_1, \ \Delta u = g_2, & \text{ on } \partial\Omega; \\ v = \psi_1, \ \Delta v = \psi_2 & \text{ on } \partial\Omega \end{cases}$$

with  $Q_1, Q_2 \in L^p(\Omega)$  and  $e^{4u}, e^{4v} \in L^{p'}(\Omega)$  for some  $1 , where <math>g_1, g_2, \psi_1, \psi_2 \in L^{\infty}(\partial\Omega)$  and  $f, g \in L^q(\Omega)$  for some q > 1. Then  $u, v \in L^{\infty}(\Omega)$ .

It follows from Theorem 2.2 that, for any solution (u, v) of system (\*),  $\int_{\Omega} Q_1(x)e^{4v}dx < \infty$ ,  $\int_{\Omega} Q_2(x)e^{4u}dx < \infty$ . Our next result states that there is a uniform bound for those integrals. For that matter, due to the fact that we are considering non-autonomous problems, we need in the theorems below geometric assumptions concerning the behavior of  $Q_1$  and  $Q_2$  near the boundary. So,

 $(H_1)$  There exist  $r, \delta > 0$  such that  $Q_1(x), Q_2(x) \in C^1(\Omega_r)$ , and

$$\nabla Q_1(x) \cdot \theta \leq 0 \text{ and } \nabla Q_2(x) \cdot \theta \leq 0$$

for all  $x \in \Omega_r(\Omega_r) := \{x \in \overline{\Omega} : dist(x, \partial \overline{\Omega}) \leq r\}$ , and unit vectors  $\theta$  such that  $|\theta - \nu(\overline{x})| < \delta$ , where  $\overline{x}$  is the closest point to x in  $\partial\Omega$  and  $\nu(\overline{x})$  denotes the unit external normal to  $\partial\Omega$  in the point  $\overline{x}$ .

With assumption  $(H_1)$  one can use the Moving Planes techniques to get bounds for the functions u and v near the boundary.

Let  $\phi_1$  be the eigenfunction associated to the first eigenvalue  $\lambda_1$  of  $(\Delta^2, H^2(\Omega) \cap H^1_0(\Omega))$ .

**Theorem 2.3.** Assume  $Q_i(x)$ , i = 1, 2 is continuous function with  $m_i \leq Q_i \leq M_i$ for some positive constants  $m_i$  and  $M_i$  and  $(H_1)$ . Assume furthermore that  $\Omega$  is convex. Then there exists a positive constant C, depending only on  $Q_i$ , i = 1, 2 and  $\Omega$ , such that

$$\int_{\Omega} Q_1(x) e^{4v} dx < C, \quad \int_{\Omega} Q_2(x) e^{4u} dx < C$$

$$(2.2)$$

for all (u, v) solution of (\*).

**Proof.** Step 1 For each (u, v) solution of system (\*) we have

$$\int_{\Omega} Q_1(x) e^{4v} \phi_1 dx \le C, \ \int_{\Omega} Q_2(x) e^{4u} \varphi_1 dx \le C,$$

where the constant C depends only on  $Q_1, Q_2$  and  $\Omega$ .

From our basic assumptions for  $Q_i$ , i = 1, 2, we know that there are positive constants  $a_i$ , i = 1, 2 with  $a_1 \times a_2 > \lambda_1^2$  and c such that

$$Q_1(x)e^{4t} \ge a_1t - c \text{ and } Q_2(x)e^{4t} \ge a_2t - c.$$
 (2.3)

Next, multiplying the equations in (\*) by  $\phi_1$ , integrating by parts and using (2.3), we obtain

$$\int_{\Omega} Q_1(x)e^{4v}\phi_1 dx = \lambda_1 \int_{\Omega} u\phi_1 dx \ge a_1 \int_{\Omega} v\phi_1 dx - c_1$$
$$\int_{\Omega} Q_2(x)e^{4u}\phi_1 dx = \lambda_1 \int_{\Omega} v\phi_1 dx \ge a_2 \int_{\Omega} u\phi_1 dx - c_1.$$
(2.4)

Thus

$$\lambda_1 \int_{\Omega} u\phi_1 dx \ge \frac{a_1 a_2}{\lambda_1} \int_{\Omega} u\phi_1 dx - c_1$$

which implies

$$\int_{\Omega} u\phi_1 dx \le C,$$

and therefore,

$$\int_{\Omega} Q_1(x) e^{4v} \phi_1 dx \le C.$$

The other inequality in (2.2) is obtained in a similar way.

Step 2 We claim that there exist  $r, \delta > 0$  such that

$$\nabla u(x) \cdot \theta \leq 0$$
 and  $\nabla v(x) \cdot \theta \leq 0$  for all  $x \in \Omega_r$ ,  $|\theta - \nu(x)| < \delta$ ,

for each (u, v) solutions of (\*), where  $\theta$  and  $\nu$  are as in  $(H_1)$ .

We can assume, without loss of generality, that  $\Omega \subset R_+^4 := \{(x_1, x_2, x_3, x_4) \in \mathbb{R}^4 : x_1 > 0\}$  and  $(0, 0, 0, 0) \in \partial \Omega$ . Now, we consider  $T_{\lambda} := \{(x_1, x_2, x_3, x_4) : x_1 = \lambda\}$ , the cap  $\Sigma_{\lambda} := \{(x_1, x_2, x_3, x_4) \in \Omega : x_1 < \lambda\}$  and the reflected cap  $\Sigma'_{\lambda} := \{(2\lambda - x_1, x_2, x_3, x_4) : (x_1, x_2, x_3, x_4) \in \Sigma_{\lambda}\}$ . It follows that there exists  $\overline{\lambda}$  such that  $\Sigma_{\lambda} \cup \Sigma'_{\lambda} \subset \Omega_r$  for each  $0 < \lambda < \overline{\lambda}$ . In fact this  $\overline{\lambda}$  depends only on r and not on the particular point on the boundary.

For  $0 < \lambda < \overline{\lambda}$ , define in  $\Sigma_{\lambda}$  the auxiliary functions

$$w_{\lambda}(x_1, x_2, x_3, x_4) = u(2\lambda - x_1, x_2, x_3, x_4) - u(x_1, x_2, x_3, x_4),$$
  
$$z_{\lambda}(x_1, x_2, x_3, x_4) = v(2\lambda - x_1, x_2, x_3, x_4) - v(x_1, x_2, x_3, x_4).$$

Using condition  $(H_1)$  we have

$$\Delta^2 w_{\lambda} = Q_1(2\lambda - x_1, x_2, x_3, x_4) e^{4v(2\lambda - x_1, x_2, x_3, x_4)} - Q_1(x_1, x_2, x_3, x_4) e^{4v(x_1, x_2, x_3, x_4)}$$

$$\geq Q_1(2\lambda - x_1, x_2, x_3, x_4) e^{4v(2\lambda - x_1, x_2, x_3, x_4)} - Q_1(2\lambda - x_1, x_2, x_3, x_4) e^{4v(x_1, x_2, x_3, x_4)}.$$

Now, using the mean value theorem we see that

$$\Delta^2 w_{\lambda} \ge c(x_1, x_2, x_3, x_4)(v(2\lambda - x_1, x_2, x_3, x_4) - v(x_1, x_2, x_3, x_4)),$$

where

$$c(x_1, x_2, x_3, x_4) = 4Q_1(2\lambda - x_1, x_2, x_3, x_4)e^{4\eta(x_1, x_2, x_3, x_4)} \ge 0$$

and  $\eta(x_1, x_2, x_3, x_4)$  is real number between  $v(x_1, x_2, x_3, x_4)$  and  $v(2\lambda - x_1, x_2, x_3, x_4)$ . Thus

$$\Delta^2 w_{\lambda} - c(x_1, x_2, x_3, x_4) z_{\lambda} \ge 0.$$

Similarly we can prove that

$$\Delta^2 z_{\lambda} - \bar{c}(x_1, x_2, x_3, x_4) w_{\lambda}(x_1, x_2, x_3, x_4) \ge 0,$$

where

$$\bar{c}(x_1, x_2, x_3, x_4) = 4Q_2(2\lambda - x_1, x_2, x_3, x_4)e^{4\xi(x_1, x_2, x_3, x_4)} \ge 0$$

and  $\xi(x_1, x_2, x_3, x_4)$  is real number between  $u(x_1, x_2, x_3, x_4)$  and  $u(2\lambda - x_1, x_2, x_3, x_4)$ .

For  $\lambda$  sufficiently small and positive we have that  $\Sigma_{\lambda}$  has small measure and so we can use the maximum principle for cooperative elliptic systems in small domains (see [9, 10]) to conclude that

$$w_{\lambda} \geq 0$$
 and  $z_{\lambda} \geq 0$  in  $\Sigma_{\lambda}$ .

Using similar arguments as in [10] we can also prove that

$$w_{\lambda} \geq 0$$
 and  $z_{\lambda} \geq 0$  in  $\Sigma_{\bar{\lambda}}$ .

Therefore, there exists  $\epsilon > 0$  such that u and v are increasing in  $\Omega_{\epsilon}$ . Finally, the conclusion follows in a standard way as in [11].

Step 3 We claim that there exist  $\epsilon > 0$  and C > 0 which depend only on  $Q_1, Q_2$ and  $\Omega$  such that  $||u||_{L^{\infty}(\Omega_{\epsilon})}, ||v||_{L^{\infty}(\Omega_{\epsilon})} \leq C$ , for each (u, v) solution of (\*).

The conclusion follows by the same arguments as in [11], using Step 2 above. Step 4 We claim that our theorem holds.

Let  $\alpha := \inf \{ \phi_1(x) : x \in \overline{\Omega} \setminus \Omega_{\epsilon} \}$ . Using Step 3 we obtain that  $Q_1(x)e^{4v}$  is bounded in  $\Omega_{\epsilon}$ . Thus

$$\int_{\Omega} Q_1(x) e^{4v} dx = \int_{\Omega_{\epsilon}} Q_1(x) e^{4v} dx + \int_{\Omega \setminus \Omega_{\epsilon}} Q_1(x) e^{4v} dx$$
$$\leq C + \frac{1}{\alpha} \int_{\Omega \setminus \Omega_{\epsilon}} Q_1(x) e^{4v} \phi_1 dx$$
$$\leq C,$$

where we have used Step1 to estimate the last integral. Using a similar argument we can prove the result for  $Q_2(x)e^{4u}$ .

Now, we presents our main result in this section.

**Theorem 2.4.** Assume  $Q_i(x)$ , i = 1, 2 is continuous function with  $m_i \leq Q_i \leq M_i$ for some positive constants  $m_i$  and  $M_i$  and  $(H_1)$ . Assume furthermore that  $\Omega$  is convex. Then there exists a constant C > 0 such that

$$||u||_{L^{\infty}}$$
 and  $||v||_{L^{\infty}} \leq C$ ,

for all eventual solutions (u, v) of system (\*).

**Proof.** In view of  $\int_{\Omega} Q_1(x)e^{4v}dx < C$  and  $\int_{\Omega} Q_2(x)e^{4u}dx < C$ , we may assume that there exist two nonnegative bounded measures  $\mu$  and  $\nu$  such that

$$Q_1(x)e^{4v_n} \to \mu \text{ and } Q_2(x)e^{4u_n} \to \nu.$$
 (2.5)

We also observe that, as a consequence of Theorem 2.3, the solutions  $((u_n, v_n))$  of (\*) are bounded in  $L^1(\Omega)$ :

$$||u_n||_{L^1}, \ ||v_n||_{L^1} \le C, \ \forall n.$$
(2.6)

A point  $x \in \Omega$  is called a  $8\pi^2$  regular point with respect to  $\mu$  if there is a function  $\psi \in C_c(\Omega), \ 0 \le \psi \le 1$ , with  $\psi = 1$  in a neighborhood of x such that

$$\int_{\Omega} \psi d\mu < 8\pi^2.$$

We define

$$\Omega_{\mu}(8\pi^2) = \{x \in \Omega : x \text{ is not a } 8\pi^2 \text{ regular point with respect to}\mu\}$$

Since  $\int d\mu < c$ , it follows that  $\Omega_{\mu}(8\pi^2)$  is a finite set. Let  $S_u$  be the blow-up set for the sequence  $(u_n)$ , that is

$$S_u := \{ x \in \Omega : \exists (x_n) \subset \Omega \text{ such that } x_n \to x \text{ and } u_n(x_n) \to +\infty \}.$$

In fact, our theorem will be proved if we can show that  $S_u = S_v = \emptyset$ .

Next, we prove our above conclusion by four steps.

Step 1 We claim that for  $x_0$  is a regular point for the measure  $\mu$  (or for the measure  $\nu$ ), then there exist constants  $\rho > 0$  and C, independent of n, such that

$$||u_n||_{L^{\infty}(B_{\rho}(x_0))} \le C, ||v_n||_{L^{\infty}(B_{\rho}(x_0))} \le C.$$

Using the fact that  $x_0$  is a regular point of the measure  $\mu$  we have a function  $\psi \in C_c(\Omega)$ ,  $0 \leq \psi \leq 1$ , with  $\psi \equiv 1$  in some neighborhood  $V_{x_0}$  of  $x_0$ , such that  $\int \psi d\mu < 8\pi^2$ . Thus,  $\int_{V_{x_0}} d\mu < 8\pi^2$ , which implies that there exist R > 0,  $\delta > 0$  and  $n_0$  such that for all  $n \geq n_0$ 

$$\int_{B_R(x_0)} Q_2(x) e^{4u_n} \le 8\pi^2 - \delta.$$
(2.7)

Using this estimates, we first work with the second equation in (\*). Let us write  $v_n := v_{1,n} + v_{2,n}$ , where

$$\triangle^2 v_{1,n} = Q_2(x)e^{4u_n}$$
, in  $B_R(x_0)$  and  $v_{1,n} = \triangle v_{1,n} = 0$  on  $\partial B_R(x_0)$ .

Notice that  $\triangle^2 v_{2,n} = 0$  in  $B_R(x_0)$ .

Using Lemma 2.1 and (2.7), we obtain

$$C \ge \int_{B_R} e^{(32\pi^2 - \frac{\delta}{2}) \frac{v_{1,n}}{\int Q_2(x)e^{4u_n}}} \ge \int_{B_R} e^{p4v_{1,n}}, \qquad (2.8)$$

where p > 1 is a constant depending only on  $\delta$ . It follows from  $t < e^t$  that

$$||v_{1,n}||_{L^p(B_R(x_0))} \le C.$$
(2.9)

Since the function  $\Delta v_{2,n}$  is harmonic, we have

$$||v_{2,n}||_{L^1(B_R)} \le ||v_n||_{L^1(B_R)} + ||v_{1,n}||_{L^1(B_R)} \le C,$$

and so

$$||v_{2,n}||_{L^{\infty}(B_{\frac{R}{4}})} \le C.$$
(2.10)

On the other hand, from  $Q_1(x)e^{4v_n} \leq ce^{4v_{1,n}}e^{4v_{2,n}}$ , we have

$$||Q_1(x)e^{4v_n}||_{L^p(B_{\frac{R}{4}})} \le c||e^{4v_n}||_{L^p(B_{\frac{R}{4}})} \le C, \text{ for some } p > 1.$$
(2.11)

In order to prove that  $||v_n||_{L^{\infty}(B_{\rho})} \leq C$ , for some  $\rho < \frac{R}{4}$ , it is enough to prove a similar bounds as (2.10) for  $v_{1,n}$ , namely

$$||v_{1,n}||_{L^{\infty}(B_{\rho})} \le C.$$
(2.12)

For that matter, we use the first equation in (\*). Let us write  $u_n = u_{1,n} + u_{2,n}$ , where

$$\triangle^2 u_{1,n} = Q_1(x)e^{4v_n}$$
, in  $B_{\frac{R}{4}}$  and  $u_{1,n} = \triangle u_{1,n} = 0$  on  $\partial B_{\frac{R}{4}}$ 

Observe that in view of (2.11), by standard elliptic regularity we have

$$||u_{1,n}||_{L^{\infty}(B_{\frac{R}{3}})} \le C.$$
(2.13)

Notice that  $\triangle^2 u_{2,n} = 0$  in  $B_{\frac{R}{4}}$ . Thus  $\triangle u_{2,n}$  is harmonic in  $B_{\frac{R}{4}}$ , and it follows that

$$||u_{2,n}||_{L^{\infty}(B_{\frac{R}{16}})} \le C.$$
(2.14)

From (2.13) and (2.14) we have

$$||u_n||_{L^{\infty}(B_{\frac{R}{16}})} \le C.$$
(2.15)

Now we go back to the second equation in (\*). Using (2.15) and elliptic regularity we have

$$||v_{2,n}||_{L^{\infty}(B_{\frac{R}{16}})} \le C.$$
(2.16)

From (2.12) and (2.16) we have

$$||v_n||_{L^{\infty}(B_{\frac{R}{16}})} \le C,$$

which together with (2.15) proves our Step 1, taking  $\rho = \frac{R}{16}$ .

Step 2 We claim that  $S_{\mu} \subset \Omega_{\mu}$  and  $S_{\nu} \subset \Omega_{\nu}$ .

In fact, This follows directly from Step 1 and the definition of the sets  $\Omega_{\mu}$ ,  $S_{\mu}$ ,  $S_{\nu}$  and  $\Omega_{\nu}$ .

Step 3 We claim that  $\Omega_{\mu} \subset S_{\nu}$  and  $\Omega_{\nu} \subset S_{\mu}$ .

Let  $x_0 \in \Omega_{\mu}$ . We claim that for each R > 0 we have

$$\lim_{n \to +\infty} ||u_n||_{L^{\infty}(B_R(x_0))} = +\infty.$$
(2.17)

Suppose by contradiction that there exists  $R_0 > 0$  and a subsequence, which we denote also by  $(u_n)$ , such that

$$||u_n||_{L^{\infty}(B_{R_0}(x_0))} \le C$$

So,

$$||Q_2(x)e^{4u_n}||_{L^{\infty}(B_{R_0}(x_0))} \le C,$$

which implies that for  $R < R_0$  we have

$$\int_{B_R(x_0)} Q_2(x) e^{4u_n} \le CR^4.$$

Thus, there exists  $R_1 > 0$ , such that

$$\int_{B_{R_1}(x_0)} Q_2(x) e^{4u_n} \le 8\pi^2$$

This implies that  $x_0$  is a regular point of  $\mu$ , which is a contradiction.

Now we observe that there exists R > 0 such that  $x_0$  is the only non-regular point in  $B_R(x_0)$ .

Next, we use (2.17) to prove that  $x_0 \in S_{\nu}$ . Indeed, from (2.17) there exists  $(x_n) \subset B_R(x_0)$  such that  $x_n \to \tilde{x}$  and  $v(x_n) \to +\infty$ . So, one needs to prove  $\tilde{x} = x_0$ . Indeed if this were not the case, then  $\tilde{x}$  would be a regular point, which is not possible, since  $u_n$  is bounded in a neighborhood of a regular point.

With similar arguments as in the proof we just completed, we can prove that  $\Omega_{\nu} \subset S_{\mu}$ .

As a consequence of Step 2 and Step 3 we conclude that those four sets coincide:

$$S_{\mu} = \Omega_{\mu} = S_{\nu} = \Omega_{\nu}.$$

Step 4 We claim that  $S_{\mu} = \emptyset$ .

We prove this claim by contradiction. Suppose that there exist  $x_0 \in S_{\mu}$ . Since  $x_0$  is isolated, we can take R > 0 such that  $\overline{B_R(x_0)} \cap (S_{\mu} \setminus \{x_0\}) = \emptyset$ .

Next, we consider the Navier boundary value problems in  $B_R(x_0)$ ,

$$\triangle^2 z_n = Q_2(x)e^{4u_n}$$
, in  $B_R(x_0)$  and  $z_n = \triangle z_n = 0$  on  $\partial B_R(x_0)$ .

We know that the function  $u_n$  satisfies

$$\Delta^2 v_n = Q_2(x)e^{4u_n}$$
, in  $B_R(x_0)$  and  $v_n \ge 0, \Delta v_n \le 0$  on  $\partial B_R(x_0)$ .

Thus, by the maximum principle we have

$$0 \le z_n \le v_n$$
 in  $\overline{B_R(x_0)}$ .

Taking the limit we have that  $z_n \to z$ , where z is a solution of the problem

$$\triangle^2 z = \mu$$
, in  $B_R(x_0)$  and  $z = \triangle z = 0$  on  $\partial B_R(x_0)$ 

On the other hand the problem

$$\triangle^2 w = 8\pi^2 \delta_0$$
, in  $B_R(x_0)$  and  $w = \triangle w = 0$  on  $\partial B_R(x_0)$ 

has the solution

$$w(x) = \ln \frac{R}{|x - x_0|}$$

Since  $x_0$  is not a regular point it follows that  $\mu > 8\pi^2 \delta_0$ . So

$$z(x) \ge ln|x - x_0|^{-1} + o(1), \ x \to x_0.$$

Now with the hypothesis  $Q_1(x)e^{4t} \ge Ce^{4t}$ , we have

$$\lim_{n \to +\infty} \int_{B_R(x_0)} Q_1(x) e^{4v_n} \ge C \int_{B_R(x_0)} e^{4w} = \infty,$$

which is impossible.

### **3.** Qualitative properties of solutions of problem (\*\*)

In this section, we study the qualitative properties of solutions of problem (\*\*).

From [8], Brezis-Merle implies that u is bounded from above when u satisfies  $-\Delta u = V(x)e^u$  and other conditions. This result is used to study the qualitative properties and classification of solutions for some second order elliptic equation (See [12, 13]). Now, one naturally ask: is any solution (u, v) to system (\*\*) with  $\int_{R^4} Q_1(x)e^{4v} < +\infty$  and  $\int_{R^4} Q_2(x)e^{4u} < +\infty$  bounded from above? We will partially answer this problem and obtain the following result:

**Theorem 3.1.** Assume  $Q_i(x)$ , i = 1, 2 is a positive bounded away from 0 and bounded from above function and (u, v) is a  $C^2$  solution of (\*\*) with  $\int_{R^4} e^{4u} < +\infty$ ,  $u(x) = o(|x|^2)$  and  $\int_{R^4} e^{4v} < +\infty$ ,  $v(x) = o(|x|^2)$  Then  $u^+ \in L^{\infty}(R^4)$  and  $v^+ \in L^{\infty}(R^4)$ .

Before we begin our proof, we need following lemmas:

**Lemma 3.1.** [3, 5] Suppose (u, v) is a  $C^2$  function on  $\mathbb{R}^4$  such that

(a)  $Q_1 e^{4v}$  and  $Q_2 e^{4u}$  are in  $L^1(\mathbb{R}^4)$  with  $0 < m_i \le Q_i \le M_i$ , i = 1, 2 for some constants  $m_i, M_i$ ;

(b) in the sense of weak derivative, u, v respectively satisfies the following equations:

$$\Delta u + \frac{2}{\beta_0} \int_{R^4} \frac{Q_1(y) e^{4v(y)}}{|x - y|^2} dy = 0$$

and

$$\Delta v + \frac{2}{\beta_0} \int_{R^4} \frac{Q_2(y)e^{4u(y)}}{|x-y|^2} dy = 0.$$

Then there are two constants  $c_1, c_2 > 0$ , respectively depending on v, u, such that  $|\Delta u|(x) \leq c_1$  on  $\mathbb{R}^4$  and  $|\Delta v|(x) \leq c_2$  on  $\mathbb{R}^4$ . Where  $\beta_0$  being given by  $(-\Delta_x)^2 (\ln \frac{1}{|x-y|}) = \beta_0 \delta_y(x)$ . In fact,  $\beta_0 = 8\pi^2$ .

**Lemma 3.2.** [5] Suppose S is  $C^2$  function on  $R^4$  such that  $0 \leq (-\Delta)S(x) \leq A$  on  $R^4$  for some constant A and  $\int_{R^4} Q(y)e^{4S(y)}dy = \alpha < \infty$  with  $0 < m \leq Q \leq M$ . Then there exists a constant B, depending only on A, m, M and  $\alpha$  such that  $S(x) \leq B$  on  $R^4$ .

**Lemma 3.3.** Suppose (u, v) is a solution of (\*\*). Let

$$w_1(x) = \frac{1}{8\pi^2} \int_{R^4} \ln \frac{|x-y|}{|y|+1} Q_1(y) e^{4v(y)} dy$$

and

$$w_2(x) = \frac{1}{8\pi^2} \int_{R^4} \ln \frac{|x-y|}{|y|+1} Q_2(y) e^{4u(y)} dy.$$

Then there exist two constants  $c_1, c_2$  such that

$$w_1(x) \le \beta_1 ln(|x|+1) + c_1$$

and

$$w_2(x) \leq \beta_2 ln(|x|+1) + c_2,$$
  
where  $\beta_1 = \frac{1}{8\pi^2} (\int_{R^4} Q_1(y) e^{4v(y)} dy)$  and  $\beta_2 = \frac{1}{8\pi^2} (\int_{R^4} Q_2(y) e^{4u(y)} dy).$ 

**Proof.** For  $|x| \ge 4$ , we decompose  $R^4 = A_1 \cup A_2$ , where  $A_1 = \{y | |y - x| \le \frac{|x|}{2}\}$ and  $A_2 = \{y | |y - x| \ge \frac{|x|}{2}\}$ . For  $y \in A_1$ , we have  $|y| \ge |x| - |x - y| \ge \frac{|x|}{2} \ge |x - y|$ , which implies

$$\ln\frac{|x-y|}{|y|+1} \le 0.$$

Since  $|x - y| \le |x| + |y| \le |x|(|y| + 1)$  for  $|x|, |y| \ge 2$  and  $\ln|x - y| \le \ln|x| + c$  for

 $|x| \ge 4$  and  $|y| \le 2$ , we have

$$w_{1}(x) \leq \frac{1}{8\pi^{2}} \int_{A_{2}} ln \frac{|x-y|}{|y|+1} Q_{1}(y) e^{4v(y)} dy$$
  
$$\leq \frac{1}{8\pi^{2}} (\int_{R^{4}} Q_{1}(y) e^{4v(y)} dy) ln |x| + c_{1}$$
  
$$= \beta_{1} ln (|x|+1) + c_{1}.$$

Similarly, we have

$$w_2(x) \le \beta_2 ln(|x|+1) + c_2.$$

**Lemma 3.4.** Suppose (u, v) is a solution of (\*\*) with  $u(x) = o(|x|^2)$  and  $v(x) = o(|x|^2)$ . Then  $\Delta u(x)$  and  $\Delta v(x)$  can be represented by

$$\Delta u(x) = -\frac{1}{4\pi^2} \int_{R^4} \frac{Q_1(y)e^{4v(y)}}{|x-y|^2} dy$$
(3.1)

and

$$\Delta v(x) = -\frac{1}{4\pi^2} \int_{R^4} \frac{Q_2(y)e^{4u(y)}}{|x-y|^2} dy.$$
(3.2)

**Proof.** Let  $k = u + w_1$ . It is obvious that  $\triangle^2 k \equiv 0$  in  $\mathbb{R}^4$ . Similar proof of Lin [1], we have for any  $x_0 \in \mathbb{R}^4$  and r > 0

$$2\pi^2 r^3 \exp(\frac{r^2}{2} \Delta k(x_0)) \le e^{-4k(x_0)} \int_{|x-x_0|=r} e^{4k} d\sigma.$$

Since  $k = u + w_1 \le u(x) + \beta ln|x| + c$  follows from Lemma 3.3, we have

$$r^{3-4\beta}\exp(\frac{\bigtriangleup k(x_0)}{2}r^2) \in L^1[1,+\infty].$$

Thus  $\Delta k(x_0) \leq 0$  for all  $x_0 \in \mathbb{R}^4$ . By Liouville's theorem,  $\Delta k(x) \equiv -c_1$  in  $\mathbb{R}^4$  for some constant  $c_1 \geq 0$ . Hence, we have

$$\Delta u(x) = -\frac{1}{4\pi^2} \int_{R^4} \frac{Q_1(y)e^{4v(y)}}{|x-y|^2} dy - c_1.$$
(3.3)

Now, we claim that  $c_1 = 0$ . Otherwise, we have  $\Delta u(x) \leq -c_1 < 0$  for  $|x| \geq R_0$ where  $R_0$  is sufficiently large. Let

$$h(y) = u(y) + \epsilon |y|^2 + A(|y|^{-2} - R_0^{-2}), \qquad (3.4)$$

where  $\epsilon$  is small such that

$$\Delta h(y) = \Delta u + 8\epsilon < -\frac{c_1}{2} < 0 \tag{3.5}$$

for  $|y| > R_0$ , and A is sufficiently large so that  $\inf_{\substack{|y| \ge R_0}} h(y)$  is achieved by some  $y_0 \in R^4$  with  $|y_0| > R_0$ . Applying the maximum principle to (3.5) at  $y_0$ , we have a contradiction. Hence, our claim is proved.

Similarly, we can prove that (3.2) holds.

**Proof of Theorem 3.1.** By Lemma 3.2 and Lemma 3.4, our conclusion holds.

Now, we study the qualitative properties of solutions of equation (\*\*). Following our Theorem 3.1 and Chen [13], we obtain the following results:

**Theorem 3.2.** Assume that  $Q_i(x), i = 1, 2$  is a positive  $C^1$  function bounded away from 0 and from above and (u, v) is a  $C^2$  solution of equation (\*\*) with  $\int_{R^4} e^{4u} dx < \infty$  and  $\int_{R^4} e^{4v} dx < \infty$ ,  $u(x) = o(|x|^2)$  and  $v(x) = o(|x|^2)$ . Then

$$-\beta_1 ln(|x|+1) - c \le u(x) \le -\beta_1 ln(|x|+1) + c \tag{3.6}$$

with  $\beta_1 > 1$  and

$$-\beta_2 ln(|x|+1) - c \le v(x) \le -\beta_2 ln(|x|+1) + c \tag{3.7}$$

with  $\beta_2 > 1$ .

Furthermore, we have the following identity

$$\int_{R^4} [(x, \nabla Q_1)e^{4v} + (x, \nabla Q_2)e^{4u}]dx = 32\pi^2 [\beta_1\beta_2 - (\beta_1 + \beta_2)].$$
(3.8)

**Theorem 3.3.** Suppose (u, v) satisfies the assumptions of Theorem 3.2 and  $Q_i$ , i = 1, 2 is radially symmetric and monotone decreasing, then u and v are radially symmetric and monotone decreasing.

**Lemma 3.5.** Assume (u, v) satisfies the assumptions of Theorem 3.2, then

$$\frac{w_i(x)}{\ln|x|} \to \beta_i, \text{ uniformly as } |x| \to \infty.$$

**Proof.** Here we prove  $w_1(x) \to \beta_1 \ln |x|$  as  $|x| \to \infty$ . We need only to verify that

$$I = \int_{R^4} \frac{\ln|x-y| - \ln(|y|+1) - \ln|x|}{\ln|x|} Q_1(y) e^{4v(y)} dy \to 0 \text{ as } |x| \to \infty.$$

Write  $I = I_1 + I_2 + I_3$ , are the integrals on the regions  $D_1 = \{y : |x - y| \le 1\}$ ,  $D_2 = \{y : |x - y| > 1 \text{ and } |y| \le k\}$  and  $D_3 = \{y : |x - y| > 1 \text{ and } |y| > k\}$  respectively. We may assume that  $|x| \ge 3$ .

(a) To estimate  $I_1$ , we simply notice that

$$|I_1| \le C \int_{|x-y|\le 1} Q_1(y) e^{4v(y)} dy - \frac{1}{\ln|x|} \int_{|x-y|\le 1} \ln|x-y| Q_1(y) e^{4v(y)} dy.$$

Then by the boundedness of  $Qe^{4v}$  (See Theorem 3.1) and  $\int_{\mathbb{R}^4} Q_1(y)e^{4v(y)}dy$ , we see that  $I_1 \to 0$  as  $|x| \to \infty$ .

(b) For each fixed k, in region  $D_2$ , we have, as  $|x| \to \infty$ ,

$$\frac{\ln|x-y| - \ln(|y|+1) - \ln|x|}{\ln|x|} \to 0.$$

Hence  $I_2 \rightarrow 0$ .

(c) To see  $I_3 \to 0$ , we use the fact that for |x - y| > 1

$$|\frac{ln|x-y| - ln(|y|+1) - ln|x|}{ln|x|}| \le c.$$

Then let  $k \to \infty$ .

Similarly, we have  $w_2(x) \to \beta_2 ln |x|$  as  $|x| \to \infty$ .

**Lemma 3.6.** Assume (u, v) satisfies the assumptions of Theorem 3.2, then

$$u(x) = \frac{1}{8\pi^2} \int_{R^4} \ln \frac{|y|+1}{|x-y|} Q_1(y) e^{4v(y)} dy + c_0$$

and

$$v(x) = \frac{1}{8\pi^2} \int_{R^4} ln \frac{|y|+1}{|x-y|} Q_2(y) e^{4u(y)} dy + \tilde{c}_0$$

where  $c_0$  and  $\tilde{c_0}$  are two constants.

**Proof.** By the Lemma 3.4, we have  $\triangle(u+w_1) = 0$  in  $\mathbb{R}^4$ . By Theorem 3.1, we have  $u^+ \in L^{\infty}$ . So, combing lemma 3.3, we have  $u + w_1 \leq cln|x| + c$ , since  $u + w_1$  is harmonic function, by the gradient estimates of harmonic functions, we have  $u(x) + w_1(x) \equiv c_0$ . Similarly, we have  $v(x) + w_2(x) \equiv \tilde{c_0}$ .

**Lemma 3.7.** Suppose (u, v) satisfies the assumptions of Theorem 3.2, then  $u_1(x) \ge -\beta_1 ln(|x|+1) - c_1$  with  $\beta_1 > 1$  and  $u_2(x) \ge -\beta_2 ln(|x|+1) - c_2$  with  $\beta_2 > 1$ .

**Proof.** By Lemma 3.3 and Lemma 3.6, we have

$$u(x) > -\beta_1 ln(|x|+1) - c_1$$

and

$$v(x) > -\beta_2 ln(|x|+1) - c_2.$$

From above inequality,  $\int_{R^4} e^{4v} dx < +\infty$  and  $\int_{R^4} e^{4u} dx < +\infty$ , we have  $\beta_1 > 1, \beta_2 > 1$ .

**Lemma 3.8.** Suppose (u, v) satisfies the assumptions of Theorem 3.2, then  $u(x) \leq -\beta_1 ln(|x|+1) + c_1$  and  $v(x) \leq -\beta_2 ln(|x|+1) + c_2$ .

**Proof.** In fact, for  $|x - y| \ge 1$ , we have

$$|x| \le |x - y|(|y| + 1)$$

Then

$$ln|x| - 2ln(|y| + 1) \le ln|x - y| - ln(|y| + 1).$$

Consequently,

$$\begin{split} w_1(x) &\geq \frac{1}{8\pi^2} \int_{|x-y| \geq 1} (\ln|x| - 2\ln(|y|+1))Q_1(y)e^{4v(y)}dy \\ &+ \frac{1}{8\pi^2} \int_{|x-y| \leq 1} (\ln|x-y| - \ln(|y|+1))Q_1(y)e^{4v(y)}dy \\ &\geq \beta_1 \ln|x| - \frac{\ln|x|}{8\pi^2} \int_{|x-y| \leq 1} Q_1(y)e^{4v(y)}dy \\ &+ \frac{1}{8\pi^2} \int_{|x-y| \leq 1} \ln|x-y|Q_1(y)e^{4v(y)}dy \\ &- \frac{1}{8\pi^2} \int_{R^4} \ln(|y|+1)Q_1(y)e^{4v(y)}dy \\ &= \beta_1 \ln|x| + I_1 + I_2 + I_3. \end{split}$$

Taking into account of the fact (see Lemma 3.5) that

$$\frac{u(x)}{\ln|x|} \to -\beta_1 \text{ and } \beta_1 > 1$$

and by the boundedness of  $Q_1(x)$ , we have

$$I_1, I_2 \to 0 \text{ as } |x| \to \infty$$

and  $I_3$  is finite. Therefore

$$w_1(x) \ge \beta_1 ln(|x|+1) - c_1.$$

By lemma 3.6, we have

$$u(x) \le -\beta_1 ln(|x|+1) + c_1.$$

Similarly, we have

$$v(x) \le -\beta_2 ln(|x|+1) + c_2$$

**Proof of Theorem 3.2.** By Lemma 3.7 and Lemma 3.8, then (3.6) and (3.7) hold. By Lin's Lemma 2.6 and lemma 2.7 in [1], we can similarly infer that (3.8) hold.

**Proof of Theorem 3.3.** By Theorem 3.2, we have  $u(x) \to -\beta_1 \ln |x|$  as  $|x| \to \infty$ , where  $\beta_1 > 1$ . Let  $\tilde{v}(x) = -\Delta u(x)$ . By revised Lin's Lemma 2.8 [1],  $\tilde{v}(x)$  has a harmonic asymptotic expansion at  $\infty$ :

$$\begin{cases} \tilde{v}(x) &= \frac{1}{|x|^2} (2\beta_1 + \sum_{j=1}^4 \frac{a_j}{|x|^2}) + \bigcirc (\frac{1}{|x|^4}), \\ \tilde{v}_{x_i} &= -\frac{4\beta_1 x_i}{|x|^4} + \bigcirc (\frac{1}{|x|^4}), \\ \tilde{v}_{x_i x_j} &= \bigcirc (\frac{1}{|x|^4}). \end{cases}$$
(3.9)

Where  $a_j$  (j = 1 to 4) are constants. Let  $\tilde{u}(x) = -\Delta v(x)$ . Similarly, we have

$$\begin{cases} \tilde{u}(x) &= \frac{1}{|x|^2} (2\beta_2 + \sum_{j=1}^4 \frac{b_j}{|x|^2}) + \bigcirc (\frac{1}{|x|^4}), \\ \tilde{u}_{x_i} &= -\frac{4\beta_2 x_i}{|x|^4} + \bigcirc (\frac{1}{|x|^4}), \\ \tilde{u}_{x_i x_j} &= \bigcirc (\frac{1}{|x|^4}). \end{cases}$$
(3.10)

Where  $b_j$  (j = 1 to 4) are constants.

Remained proof essentially equals to Lin's proof. We omit it here.

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