

Linear Transformations of N -connections in OSC^2M (II)

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Abstract. In the tangent space $T(E)$ and in its dual $T^*(E)$, two different adapted basis was introduced, in another paper with the same title [10]. In the present paper are given the connections between $\hat{K}_{\alpha\alpha}^{\gamma\beta}$ and $K_{\alpha\alpha}^{\gamma\beta}$, between the torsions: $\hat{T}_{\alpha\alpha}^{\gamma\beta}$ and $T_{\alpha\alpha}^{\gamma\beta}$ and the components of the metric tensor $\hat{g}_{\alpha\alpha}^{\beta\beta}$ and $g_{\alpha\alpha}^{\beta\beta}$.

1. Introduction

We define $E = Osc^2M$ as a $3-n$ dimensional C^∞ real manifold, in which the transformation of form (1.1) are allowed. It is formed as a tangent space of order two of the n -dimensional base manifold M . In some local chart (U, φ) some point $u \in E$ has coordinates

$$(x^a, y^{1a}, y^{2a}) = (y^{0a}, y^{1a}, y^{2a}) = (y^{aa}),$$

where $x^a = y^{0a}$ and $a, b, c, d, e, \dots = \overline{1, n}$, $\alpha, \beta, \gamma, \delta, \dots = \overline{0, 2}$.

The following abbreviations will be used

$$\partial_{\alpha\alpha} = \frac{\partial}{\partial y^{aa}}, \quad \alpha = \overline{0, 2}.$$

If in some other chart (U', φ') the point $u \in E$ has coordinates $(x^{a'}, y^{1a'}, y^{2a'})$, then in $U \cap U'$ the allowable coordinates transformations are given by

$$\begin{cases} x^{a'} = x^{a'}(x^1, x^2, \dots, x^n), \\ y^{1a'} = (\partial_a x^{a'}) y^{1a} = (\partial_{0a} y^{0a'}) y^{1a}, \\ y^{2a'} = (\partial_{0a} y^{1a'}) y^{1a} + (\partial_{1a} y^{1a'}) y^{2a}. \end{cases} \quad (1.1)$$

The natural base of $T(E)$ and of $T^*(E)$ are respectively

$$\bar{B} = \{\partial_{0a}, \partial_{1a}, \partial_{2a}\} \quad (1.2)$$

and

$$\bar{B}^* = \{dy^{0a}, dy^{1a}, dy^{2a}\}. \quad (1.3)$$

The elements of \bar{B} and \bar{B}^* are dual to each other, i.e.,

$$\langle dy^{aa}, \partial_{\beta b} \rangle = \delta_{\beta}^a \delta_b^a, \quad (1.4)$$

but with respect to (1.1) they have not a tensorial character.

The adapted basis B^* of $T^*(E)$ is given by

$$B^* = \{\delta y^{0a}, \delta y^{1a}, \delta y^{2a}\} \quad (1.5)$$

where

$$\begin{cases} \delta y^{0a} = dx^a = dy^{0a}, \\ \delta y^{1a} = dy^{1a} + M_0^1 \frac{a}{b} dy^{0b}, \\ \delta y^{2a} = dy^{2a} + M_1^2 \frac{a}{b} dy^{1b} + M_0^2 \frac{a}{b} dy^{0b}. \end{cases} \quad (1.6)$$

Theorem 1.1. [10] *The necessary and sufficient conditions that δy^{aa} are transformed as d-tensor fields are that the following equations are satisfied*

$$\begin{cases} M_0^1 \frac{a}{b} (\partial_a x^{b'}) = M_0^1 \frac{b'}{c'} (\partial_b x^{c'}) + \partial_b y^{1b'}, \\ M_1^2 \frac{a}{b} (\partial_a x^{b'}) = M_1^2 \frac{b'}{c'} (\partial_b y^{1c'}) + \partial_b y^{2b'}, \\ M_0^2 \frac{a}{b} (\partial_a x^{b'}) = M_0^2 \frac{b'}{c'} (\partial_b x^{c'}) + M_1^2 \frac{b'}{c'} (\partial_b y^{1c'}) + \partial_b y^{2b'}. \end{cases} \quad (1.7)$$

Let us denote the adapted basis of $T(E)$ by B , where

$$B = \{\delta_{0a}, \delta_{1a}, \delta_{2a}\} = \{\delta_{\alpha a}\}, \quad (\alpha = \overline{0, 2}), \quad (1.8)$$

and

$$\begin{cases} \delta_{0a} = \partial_{0a} - N_{0a}^1{}^b \partial_{1b} - N_{0a}^2{}^b \partial_{2b}, \\ \delta_{1a} = \partial_{1a} - N_{1a}^2{}^b \partial_{2b}, \\ \delta_{2a} = \partial_{2a}. \end{cases} \quad (1.9)$$

Theorem 1.2. [10] B is dual to B^* , if and only if the following relations hold

$$\begin{cases} N_{0a}^1{}^b = M_{0a}^1{}^b, \\ N_{0a}^2{}^b = M_{0a}^2{}^b - M_{1c}^2{}^b N_{0a}^1{}^c, \\ N_{1a}^2{}^b = M_{1a}^2{}^b. \end{cases} \quad (1.10)$$

Definition 1.1. *The generalized connection*

$$\nabla : T(E) \otimes T(E) \rightarrow T(E), \quad \nabla : (X, Y) \rightarrow \nabla_X Y,$$

or equivalently

$$\nabla_X : T(E) \rightarrow T(E), \quad \nabla_X : Y \rightarrow \nabla_X Y,$$

is a linear connection determined by

$$\nabla_{\delta_{\beta b}} \delta_{\alpha a} = \Gamma_{\alpha a}^{\gamma c}{}_{\beta b} \delta_{\gamma c}, \quad (1.11)$$

where the summation is going over γ and c . If in (1.11) we set $\gamma = \alpha$ this provides the so called d -connection:

$$\nabla_{\delta_{\beta b}} \delta_{\alpha a} = \Gamma_{\alpha a}^{\alpha c}{}_{\beta b} \delta_{\alpha c}, \quad (1.12)$$

(with no summation over α).

The explicit form of (1.11) is given by

$$\begin{aligned}\nabla_{\delta_{0b}} \delta_{0a} &= \underline{\Gamma_{0a}^{0c} \delta_{0c}} + \Gamma_{0a}^{1c} \delta_{1c} + \Gamma_{0a}^{2c} \delta_{2c}, \\ \nabla_{\delta_{0b}} \delta_{1a} &= \Gamma_{1a}^{0c} \delta_{0c} + \underline{\Gamma_{1a}^{1c} \delta_{1c}} + \Gamma_{1a}^{2c} \delta_{2c}, \\ \nabla_{\delta_{0b}} \delta_{2a} &= \Gamma_{2a}^{0c} \delta_{0c} + \Gamma_{2a}^{1c} \delta_{1c} + \underline{\Gamma_{2a}^{2c} \delta_{2c}}.\end{aligned}\tag{1.13}$$

If in the above formulae we substitute $0b$ with $1b$, and then $0b$ with $2b$, we obtain the complete list of 9 formulae. The underlined terms are $\Gamma_{0a}^{0c} \delta_{0c}$, $\Gamma_{1a}^{1c} \delta_{1c}$, $\Gamma_{2a}^{2c} \delta_{2c}$, where instead of X stays $1b$ or $2b$.

Comparing (1.13) with (1.12) we see that if in (1.13) all terms are zero except the underlined ones, we obtain the explicit form of the so called d-connection defined by (1.12).

Assume that $\hat{B} = \{\hat{\delta}_{0a}, \hat{\delta}_{1a}, \hat{\delta}_{2a}\}$ is another adapted basis of $T(E)$, which is formed as B(1.9) but with N replaced with \hat{N} . Another adapted basis of $T^*(E)$ is $\hat{B}^* = \{\hat{\delta}y^{0a}, \hat{\delta}y^{1a}, \hat{\delta}y^{2a}\}$ which is formed as B*(1.6) but with M replaced with \hat{M} .

2. The connection between the torsions in the new adapted basis

The torsion tensor $T(X, Y)$ is defined as usual by

$$T(X, Y) = \nabla_X Y - \nabla_Y X - [X, Y].$$

Theorem 2.1. [10] *The torsion tensor of generalized connection has the form*

$$T(X, Y) = T_{\beta b}^{\gamma c}{}_{\alpha a} X^{\alpha a} Y^{\beta b} \delta_{\gamma c},\tag{2.1}$$

where

$$\begin{aligned}T_{\beta b}^{\gamma c}{}_{\alpha a} &= \Gamma_{\beta b}^{\gamma c}{}_{\alpha a} - \Gamma_{\alpha a}^{\gamma c}{}_{\beta b} - K_{\alpha a}^{\gamma c}{}_{\beta b}, \\ [\delta_{\alpha a}, \delta_{\beta b}] &= K_{\alpha a}^{\gamma c}{}_{\beta b} \delta_{\gamma c},\end{aligned}$$

with

$$\begin{aligned}
 K_{0a}^{1c}{}_{0b} &= \delta_{0b} N_{0a}^{1c} - \delta_{0a} N_{0b}^{1c}, \\
 K_{0a}^{2c}{}_{0b} &= (\delta_{0b} N_{0a}^{2c} - \delta_{0a} N_{0b}^{2c}) + M_{1e}^2 K_{0a}^{1e}{}_{0b}, \\
 K_{0a}^{1c}{}_{1b} &= \delta_{1b} N_{0a}^{1c}, \\
 K_{0a}^{2c}{}_{1b} &= (\delta_{1b} N_{0a}^{2c} - \delta_{0a} N_{1b}^{2c}) + M_{1d}^2 K_{0a}^{1d}{}_{1b}, \\
 K_{0a}^{1c}{}_{2b} &= \delta_{2b} N_{0a}^{1c}, \\
 K_{0a}^{2c}{}_{2b} &= \delta_{2b} N_{0a}^{2c} + M_{1d}^2 K_{0a}^{1d}{}_{2b}, \\
 K_{0a}^{2c}{}_{1b} &= \delta_{1b} N_{1a}^{2c} - \delta_{1a} N_{1b}^{2c}, \\
 K_{1a}^{2c}{}_{2b} &= \delta_{2b} N_{1a}^{2c}.
 \end{aligned}$$

Using the relations (4.10), [10] we prove

Theorem 2.2. *Between $\hat{K}_{\alpha a}^{\gamma c}{}_{\beta b}$ and $K_{\alpha a}^{\gamma c}{}_{\beta b}$ there are the following connections*

$$\begin{cases}
 \hat{K}_{0a}^{0e}{}_{0b} &= 0, \\
 \hat{K}_{0a}^{1e}{}_{0b} &= K_{0b}^{1e}{}_{0a} + \mathbf{A}_{ab} \{ \bar{K}_{0a}^{1e}{}_{0b} + \hat{\delta}_{0a} A_{0b}^{1e} \}, \\
 A_{1c}^2 \hat{K}_{0a}^{1c}{}_{0b} + \hat{K}_{0a}^{2e}{}_{0b} &= K_{0b}^{2e}{}_{0a} + A_{0a}^1 A_{0b}^d K_{1c}^{2e}{}_{1d} + \mathbf{A}_{ab} \\
 &\quad \{ \bar{K}_{0a}^{2e}{}_{0b} + \hat{\delta}_{0a} \bar{A}_{0b}^2{}^e + A_{0a}^1 \bar{A}_{0b}^2{}^d K_{1c}^{2e}{}_{2d} \}.
 \end{cases} \quad (2.2)$$

$$\begin{cases}
 \hat{K}_{0a}^{0e}{}_{1b} &= 0, \\
 \hat{K}_{0a}^{1e}{}_{1b} &= \bar{K}_{0b}^{1e}{}_{1b} - \hat{\delta}_{1b} A_{0a}^{1e}, \\
 A_{1c}^2 \hat{K}_{0a}^{1c}{}_{1b} + \hat{K}_{0a}^{2e}{}_{1b} &= \bar{K}_{0a}^{2e}{}_{1b} + A_{0a}^1 \bar{K}_{1c}^{2e}{}_{1b} + \hat{\delta}_{0a} A_{1b}^{2e} - \hat{\delta}_{1b} \bar{A}_{0a}^2{}^e - \\
 &\quad - A_{0a}^2 K_{1b}^{2e}{}_{2c}.
 \end{cases} \quad (2.3)$$

$$\begin{cases}
 \hat{K}_{0a}^{0e}{}_{2b} &= 0, \\
 \hat{K}_{0a}^{1e}{}_{2b} &= K_{0a}^{1e}{}_{2b} - \delta_{2b} A_{0a}^{1e}, \\
 A_{1c}^2 \hat{K}_{0a}^{1c}{}_{2b} + \hat{K}_{0a}^{2e}{}_{2b} &= K_{0a}^{2e}{}_{2b} + A_{0a}^1 K_{1c}^{2e}{}_{2b} - \delta_{2b} \bar{A}_{0a}^2{}^e.
 \end{cases} \quad (2.4)$$

$$\begin{cases}
 \hat{K}_{1a}^{0e}{}_{1b} = 0, \\
 \hat{K}_{1a}^{1e}{}_{1b} = 0, \\
 \hat{K}_{1a}^{2e}{}_{1b} = K_{1b}^{2e}{}_{1a} + \mathbf{A}_{ab} \{ \bar{K}_{1a}^{2e}{}_{1b} + \hat{\delta}_{1a} A_{1b}^{2e} \}.
 \end{cases} \quad (2.5)$$

$$\left\{ \begin{array}{l} \hat{K}_{1a}{}^{0e}{}_{2b} = 0, \\ \hat{K}_{1a}{}^{1e}{}_{2b} = 0, \\ \hat{K}_{1a}{}^{2e}{}_{2b} = K_{1a}{}^{2e}{}_{2b} - \delta_{2b} A_1^2{}^e{}_a, \end{array} \right. \quad \left\{ \begin{array}{l} \hat{K}_{2a}{}^{0e}{}_{2b} = K_{2a}{}^{0e}{}_{2b} = 0, \\ \hat{K}_{2a}{}^{1e}{}_{2b} = K_{2a}{}^{1e}{}_{2b} = 0, \\ \hat{K}_{2a}{}^{2e}{}_{2b} = K_{2a}{}^{2e}{}_{2b} = 0, \end{array} \right. \quad (2.6)$$

where

$$\left\{ \begin{array}{l} \bar{K}_{\hat{\alpha}d}{}^{\gamma e}{}_{0b} = K_{\hat{\alpha}d}{}^{\gamma e}{}_{0b} + A_0^1{}^c{}_b K_{\hat{\alpha}d}{}^{\gamma e}{}_{1c} + \bar{A}_0^2{}^c{}_b K_{\hat{\alpha}d}{}^{\gamma e}{}_{2c}, \\ \bar{K}_{\hat{\alpha}d}{}^{\gamma e}{}_{1b} = K_{\hat{\alpha}d}{}^{\gamma e}{}_{1b} + A_1^2{}^c{}_b K_{\hat{\alpha}d}{}^{\gamma e}{}_{2c}. \end{array} \right. \quad (2.7)$$

The connections between $\hat{T}_{ca}{}^{\gamma c}{}_{\beta b}$ and $T_{ca}{}^{\gamma c}{}_{\beta b}$ are given in the following theorem.

Theorem 2.3. *The components of the same torsion tensor $T(X,Y)$, of the generalized connection ∇ in the basis B and \hat{B} are connected by*

$$\left\{ \begin{array}{l} \hat{T}_{0b}{}^{0e}{}_{0a} = \bar{T}_{0b}{}^{0e}{}_{0a} + A_0^1{}^d{}_b \bar{T}_{1d}{}^{0e}{}_{0a} + \bar{A}_0^2{}^d{}_b \bar{T}_{2d}{}^{0e}{}_{0a}, \\ \hat{T}_{0b}{}^{1e}{}_{0a} + A_0^1{}^e{}_d \hat{T}_{0b}{}^{0d}{}_{0a} = \bar{T}_{0b}{}^{1e}{}_{0a} + A_0^1{}^d{}_b \bar{T}_{1d}{}^{1e}{}_{0a} + \bar{A}_0^2{}^d{}_b \bar{T}_{2d}{}^{1e}{}_{0a}, \\ \hat{T}_{0b}{}^{2e}{}_{0a} + A_1^2{}^e{}_d \hat{T}_{0b}{}^{1d}{}_{0a} + \bar{A}_0^2{}^e{}_d \hat{T}_{0b}{}^{0d}{}_{0a} = \bar{T}_{0b}{}^{2e}{}_{0a} + A_0^1{}^d{}_b \bar{T}_{1d}{}^{2e}{}_{0a} + \bar{A}_0^2{}^d{}_b \bar{T}_{2d}{}^{2e}{}_{0a}. \end{array} \right. \quad (2.8)$$

$$\left\{ \begin{array}{l} \hat{T}_{0b}{}^{0e}{}_{1a} = \bar{T}_{0b}{}^{0e}{}_{1a} + A_0^1{}^d{}_b \bar{T}_{1d}{}^{0e}{}_{1a} + \bar{A}_0^2{}^d{}_b \bar{T}_{2d}{}^{0e}{}_{1a}, \\ \hat{T}_{0b}{}^{1e}{}_{1a} + A_0^1{}^e{}_d \hat{T}_{0b}{}^{0d}{}_{1a} = \bar{T}_{0b}{}^{1e}{}_{1a} + A_0^1{}^d{}_b \bar{T}_{1d}{}^{1e}{}_{1a} + \bar{A}_0^2{}^d{}_b \bar{T}_{2d}{}^{1e}{}_{1a}, \\ \hat{T}_{0b}{}^{2e}{}_{1a} + A_1^2{}^e{}_d \hat{T}_{0b}{}^{1d}{}_{1a} + \bar{A}_0^2{}^e{}_d \hat{T}_{0b}{}^{0d}{}_{1a} = \bar{T}_{0b}{}^{2e}{}_{1a} + A_0^1{}^d{}_b \bar{T}_{1d}{}^{2e}{}_{1a} + \bar{A}_0^2{}^d{}_b \bar{T}_{2d}{}^{2e}{}_{1a}. \end{array} \right. \quad (2.9)$$

$$\left\{ \begin{array}{l} \hat{T}_{0b}{}^{0e}{}_{2a} = T_{0b}{}^{0e}{}_{2a} + A_0^1{}^d{}_b T_{1d}{}^{0e}{}_{2a} + \bar{A}_0^2{}^d{}_b T_{2d}{}^{0e}{}_{2a}, \\ \hat{T}_{0b}{}^{1e}{}_{2a} + A_0^1{}^e{}_d \hat{T}_{0b}{}^{0d}{}_{2a} = T_{0b}{}^{1e}{}_{2a} + A_0^1{}^d{}_b T_{1d}{}^{1e}{}_{2a} + \bar{A}_0^2{}^d{}_b T_{2d}{}^{1e}{}_{2a}, \\ \hat{T}_{0b}{}^{2e}{}_{2a} + A_1^2{}^e{}_d \hat{T}_{0b}{}^{1d}{}_{2a} + \bar{A}_0^2{}^e{}_d \hat{T}_{0b}{}^{0d}{}_{2a} = T_{0b}{}^{2e}{}_{2a} + A_0^1{}^d{}_b T_{1d}{}^{2e}{}_{2a} + \bar{A}_0^2{}^d{}_b T_{2d}{}^{2e}{}_{2a}. \end{array} \right. \quad (2.10)$$

$$\left\{ \begin{array}{l} \hat{T}_{1b}{}^{0e}{}_{1a} = \bar{T}_{1b}{}^{0e}{}_{1a} + A_1^2{}^d{}_b \bar{T}_{2d}{}^{0e}{}_{1a}, \\ \hat{T}_{1b}{}^{1e}{}_{1a} + A_0^1{}^e{}_d \hat{T}_{1b}{}^{0d}{}_{1a} = \bar{T}_{1b}{}^{1e}{}_{1a} + A_1^2{}^d{}_b \bar{T}_{2d}{}^{1e}{}_{1a}, \\ \hat{T}_{1b}{}^{2e}{}_{1a} + A_1^2{}^e{}_d \hat{T}_{1b}{}^{1d}{}_{1a} + \bar{A}_0^2{}^e{}_d \hat{T}_{1b}{}^{0d}{}_{1a} = \bar{T}_{1b}{}^{2e}{}_{1a} + A_1^2{}^d{}_b \bar{T}_{2d}{}^{2e}{}_{1a}. \end{array} \right. \quad (2.11)$$

$$\begin{cases} \hat{T}_{2b\ 1a}^{0e} & = \overline{T}_{2b\ 1a}^{0e}, \\ \hat{T}_{2b\ 1a}^{1e} + A_0^1{}^e{}_d \hat{T}_{2b\ 1a}^{0d} & = \overline{T}_{2b\ 1a}^{1e}, \\ \hat{T}_{2b\ 1a}^{2e} + A_1^2{}^e{}_d \hat{T}_{2b\ 1a}^{1d} + \overline{A}_0^2{}^e{}_d \hat{T}_{2b\ 1a}^{0d} & = \overline{T}_{2b\ 1a}^{2e}. \end{cases} \quad (2.12)$$

$$\hat{T}_{2b\ 2a}^{\gamma e} = T_{2b\ 2a}^{\gamma e}, \quad \gamma = \overline{0, 2}, \quad (2.13)$$

where

$$\begin{cases} \overline{T}_{\delta d\ 0b}^{\gamma e} = T_{\delta d\ 0b}^{\gamma e} + A_0^1{}^c{}_b T_{\delta d\ 1c}^{\gamma e} + \overline{A}_0^2{}^c{}_b T_{\delta d\ 2c}^{\gamma e}, \\ \hat{T}_{\delta d\ 1b}^{\gamma e} = T_{\delta d\ 1b}^{\gamma e} + A_1^2{}^c{}_b T_{\delta d\ 2c}^{\gamma e}, \quad \delta = \overline{0, 2}. \end{cases} \quad (2.14)$$

3. The connections between the components of the metric tensor in the new adapted basis

In the space $T^*(E) \otimes T^*(E)$ the metric tensor G can be given by

$$\begin{aligned} G &= \begin{bmatrix} \delta y^{0a} & \delta y^{1a} & \delta y^{2a} \end{bmatrix} \begin{bmatrix} g_{0a\ 0b} & g_{0a\ 1b} & g_{0a\ 2b} \\ g_{1a\ 0b} & g_{1a\ 1b} & g_{1a\ 2b} \\ g_{2a\ 0b} & g_{2a\ 1b} & g_{2a\ 2b} \end{bmatrix} \begin{bmatrix} \delta y^{0b} \\ \delta y^{1b} \\ \delta y^{2b} \end{bmatrix} \\ &= g_{\alpha a\ \beta b} \delta y^{\alpha a} \otimes \delta y^{\beta b}, \quad \alpha, \beta = \overline{0, 2}. \end{aligned}$$

For the components of the metric tensor we have:

$$g_{\alpha a\ \beta b} = g_{\alpha a'\ \beta b'} (\partial_a x^{a'}) (\partial_b x^{b'}), \quad \alpha, \beta = \overline{0, 2}.$$

Theorem 3.1. *The connections between the components of the metric tensor, with respect to the generalized connection ∇ , in the basis B and \hat{B} are given by*

$$\hat{g}_{0a\ 0b} = \overline{g}_{0a\ 0b} + A_0^1{}^d{}_b \overline{g}_{0a\ 1d} + \overline{A}_0^2{}^d{}_b \overline{g}_{0a\ 2d}, \quad (3.1)$$

$$\hat{g}_{1a\ 0b} = \overline{g}_{1a\ 0b} + A_1^2{}^c{}_a \overline{g}_{2c\ 0b}, \quad (3.2)$$

$$\hat{g}_{2a\ 0b} = \overline{g}_{2a\ 0b}, \quad (3.3)$$

$$\hat{g}_{0a\ 1b} = \overline{g}_{0a\ 1b} + A_1^2{}^d{}_b \overline{g}_{0a\ 2d}, \quad (3.4)$$

$$\hat{g}_{1a\ 1b} = \overline{g}_{1a\ 1b} + A_1^2{}^d{}_b \overline{g}_{1a\ 2d}, \quad (3.5)$$

$$\hat{g}_{2a\ 1b} = \bar{g}_{2a\ 1b}, \quad (3.6)$$

$$\hat{g}_{0a\ 2b} = \bar{g}_{0a\ 2b}, \quad (3.7)$$

$$\hat{g}_{1a\ 2b} = \bar{g}_{1a\ 2b}, \quad (3.8)$$

$$\hat{g}_{2a\ 2b} = g_{2a\ 2b} \quad (3.9)$$

where

$$\bar{g}_{0a\ \gamma b} = g_{0a\ \gamma b} + A_{0\ a}^{1\ c} g_{1c\ \gamma b} + \bar{A}_{0\ a}^{2\ c} g_{2c\ \gamma b}, \quad \gamma = \overline{0,2}, \quad (3.10)$$

$$\bar{g}_{\gamma a\ 0b} = g_{\gamma a\ 0b} + A_{0\ b}^{1\ d} g_{\gamma a\ 1d} + \bar{A}_{0\ b}^{2\ d} g_{\gamma a\ 2d}, \quad \gamma = \overline{1,2}, \quad (3.11)$$

$$\bar{g}_{1a\ \gamma b} = g_{1a\ \gamma b} + A_{1\ a}^{2\ c} g_{2c\ \gamma b}, \quad (3.12)$$

$$\bar{g}_{2a\ 1b} = g_{2a\ 1b} + A_{1\ b}^{2\ d} g_{2a\ 2d}. \quad (3.13)$$

References

1. I. Čomić, Different adapted basis and metrical generalized connections, *Novi Sad J. Math.* **27**, **1** (1997), 117-132.
2. I. Čomić, Gh. Atanasiu, and E. Stoica, The generalized connections in Osc^3M , *Annales Univ. Sci. Budapest* **41** (1998), 39-54.
3. M. Matsumoto, M., *The Theory of Finsler Connections*, Publ. Study Group Geom. 5, Depart. Math., Okayama Univ., 1970.
4. M. Matsumoto, *Foundations of Finsler Geometry and Special Finsler Spaces*, Kaiseisha Press, Otsu, Japan, 1986.
5. R. Miron and M. Anastasiei, *Geometry of Lagrange Spaces, Theory and Applications*, Kluwer Academic Publishers, 1993.
6. R. Miron, *The Geometry of Higher Order Lagrange Spaces, Applications to Mechanics and Physics*, Kluwer Academic Publishers, FTPH 82, 1997.
7. R. Miron and Gh. Atanasiu, Differential geometry of the k -Osculator bundle, *Rev. Roumaine Math. Pures Appl.* **41**, **3/4** (1996), 205-236.
8. R. Miron and Gh. Atanasiu, Higher-order Lagrange spaces, *Rev. Roumaine Math. Pures Appl.* **41**, **3/4** (1996), 251-263.
9. R. Miron and M. Hashiguchi, Metrical Finsler connections, *Rep. Mat. Sci., Kagoshima Univ. (Math., Phys. & Chem.)* **12** (1979), 21-35.
10. I. Čomić and M. Purcaru, Linear transformation of N -Connection in Osc^2M , *Differential Geometry and Dynamical Systems (electronic journal)*, <http://www.mathem.pub.ro/apps/v2n1.ltm>, **2** (1) (2000), 18-31.

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